

Dyson Schwinger Eqs and Quantization of gauge theories (Summer '21)

Dirk Kreimer lect. June 21 2021

$$\phi_R(\not{q}, \not{p}_1, \not{p}_2) = \frac{\sqrt{\lambda(q^2, m_1^2, m_2^2)}}{2q^2} \left\{ \ln \left(\frac{q^2 - m_1^2 - m_2^2 + \sqrt{\lambda}}{q^2 - m_1^2 - m_2^2 - \sqrt{\lambda}} \right) + \dots \right\}$$

$$\lambda(q, k, c) = q^2 + k^2 + c^2 - 2(ak + bc + ca)$$

$$\text{Im}(\phi_R(q^2)) = \Theta(q^2 - (m_1 + m_2)^2) \frac{\sqrt{\lambda}}{2q^2}$$

Wightman axioms:

monodromy in full Green fct.

$$\text{So far } X^- = \mathbb{I} - \sum_{j=1}^{\infty} \alpha_j \sum_{i=1}^{\infty} \beta_{i,j} (X^- Q^{2j})$$

$$\phi_R(X^-)(q^2, \dots)$$

$$\text{Im} \phi_R(X^-) = \delta(q^2 - m^2) - \int_{(m_1+m_2)^2}^{\infty} \frac{S(x)}{x(x-m^2)} dx$$

Killian - representation

Wanted: $\text{Im} \phi_R(X^-)$ as a solution

To a fixed pt equation, so as a solution to a DSE.

$$\text{Diagram} = \text{Diagram} - (\text{Diagram} + \text{Diagram} + \text{Diagram} + \dots)$$

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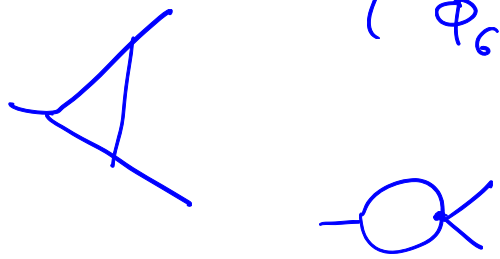
To solve that problem:

i) need core Hopf algebra

$$\Delta_{\text{core}}(\Gamma) = \Gamma \otimes \underline{1} + \underline{1} \otimes \Gamma + \sum_{\gamma \neq \Gamma} \gamma \otimes \Gamma/\gamma$$

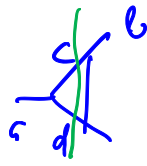
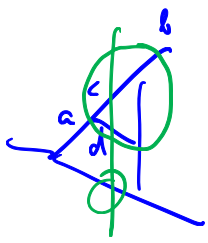
$\gamma = \parallel \gamma$

as quotient Hopf algebras, you have
 new realization Hopf algebras as needed
 monodromy of $(\phi_6^3 - \text{graph})$

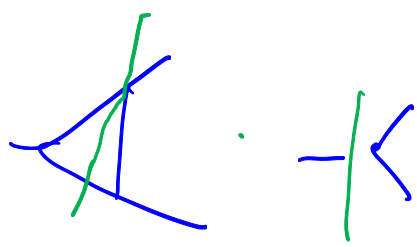


ii) Hopf algebra of We-Casimir graphs.

$$\text{Diagram} = \alpha \text{Diagram} + \alpha^2 \left(\text{Diagram} + \text{Diagram} + \text{Diagram} + \dots + \text{Diagram} + \text{Diagram} + \text{Diagram} \right)$$

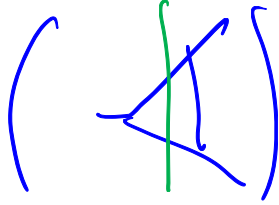
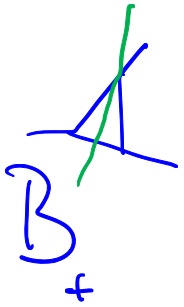


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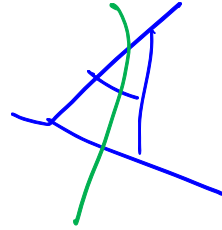


pre-cutkosky

graph



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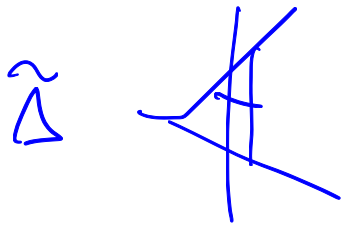


1-cocycle

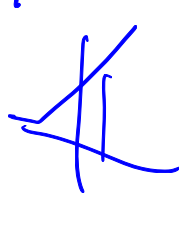
H_{pre} ?

iii) H_c vector space of

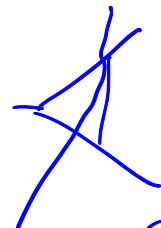
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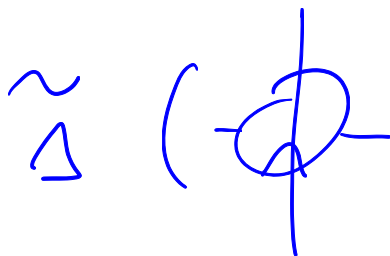


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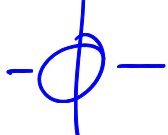
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iv) Co-actions

$$H \rightarrow H \otimes H$$

Start with us V , find a

map:

$$\bar{\Delta}: V \rightarrow H \otimes V$$

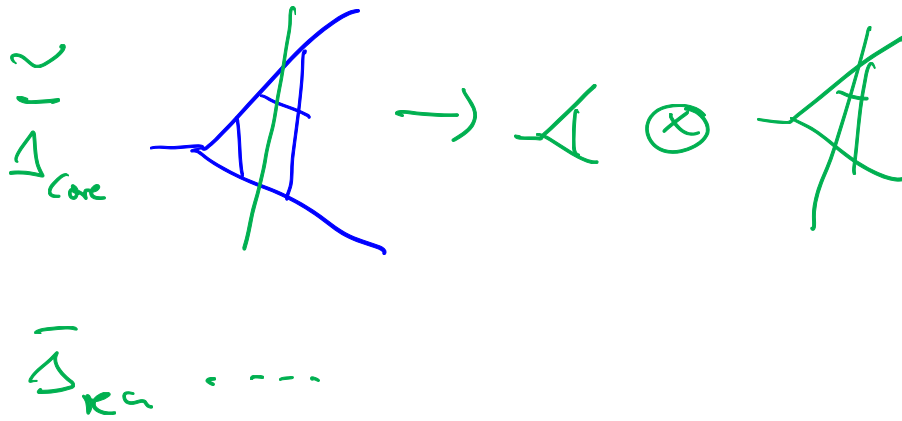
↖ Hood of Δ

such that

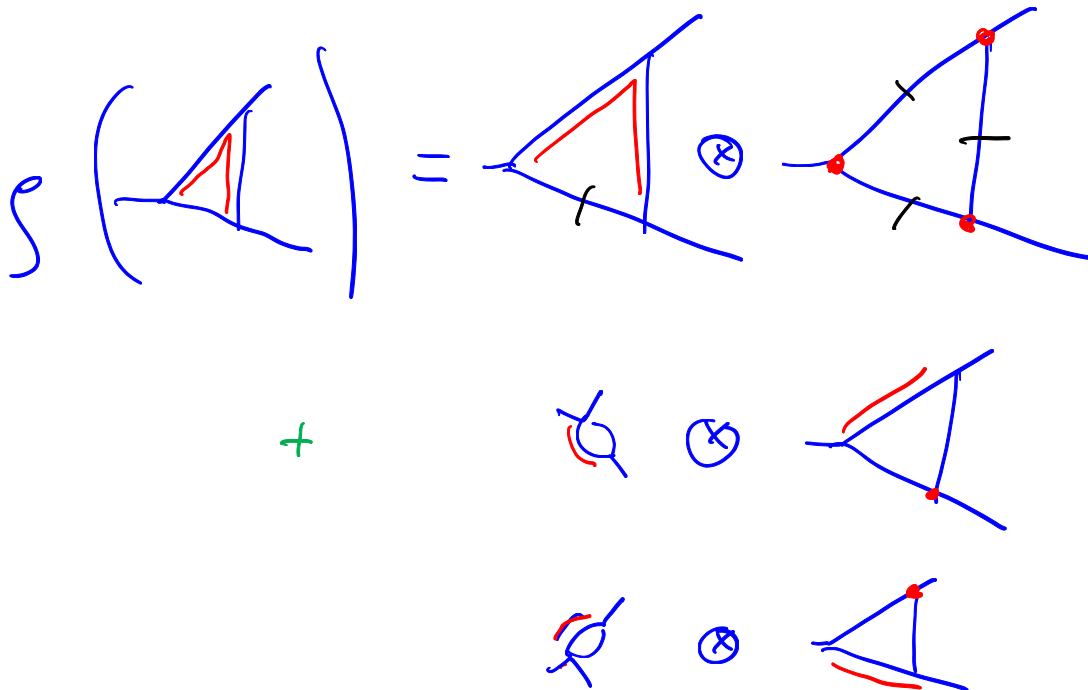
$$(id \otimes \bar{\Delta}) \bar{\Delta}: V \rightarrow H \otimes H \otimes V$$

$$(\bar{\Delta} \otimes id) \bar{\Delta}: V \rightarrow H \otimes H \otimes V$$

this gives a co-action.

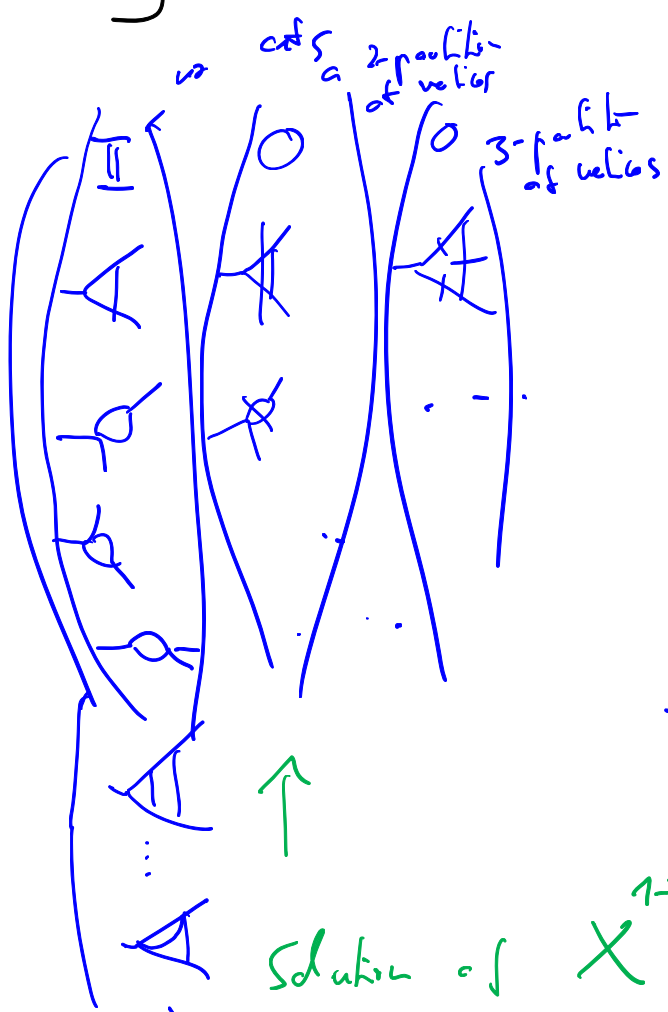


mono-dromy co-action S .



This coaction co-exists "nicely" with the one above to give a "Co-interacting Bialgebras" (work of Mascher, Foissy, ...).

↳ } a matrix Γ^{\uparrow}



these are Hodge matrices

Solution of $X = \sum B_{\pm}^{r,j} (X_{\dots})$

$X^{1,1,1}$ \uparrow \mathbb{Q}

left most column is solution of $X = \mathbb{I} \mp \sum_i a_i B_{\pm}^{r,j} (\dots)$

Lower triangular matrices which are Hodge, being Hodge \Leftrightarrow the $C^{(i,i)}$ column jins the monodromy of the column $C^{(i)}$, row by row.

Famous example: the polylog

$$Li_k(z) = \int_1^z \frac{Li_{k-1}(\tilde{z})}{\tilde{z}} d\tilde{z} \quad \forall k \geq 2$$

$$Li_2(z) = \int_1^z \frac{Li_1(\tilde{z})}{\tilde{z}} d\tilde{z} \quad Li_1(z) = -\log(1-z)$$

1	0	0	0	...
$Li_1(z)$	$2\pi i$	0	0	...
$Li_2(z)$	$2\pi i \log z$	$(2\pi i)^2$	0	...
$Li_3(z)$	$2\pi i \frac{1}{2} \log^2 z$	$(2\pi i)^2 \log z$	$(2\pi i)^3$...
\vdots	\vdots	\vdots	\vdots	\vdots
\vdots	$2\pi i \frac{1}{n!} \log^n z$	$(2\pi i)^2 \log^{n-1} z$	$(2\pi i)^3 \log^{n-2} z$...

$$z \frac{\partial Li_k(z)}{\partial z} = Li_{k-1}(z)$$

Want to do the same for Feynman graphs

