

Algebraic lattices in QFT renormalization

Michael Borinsky*

Institute of Physics, Humboldt University
Newton Str. 15, D-12489 Berlin, Germany

September 8, 2015

Abstract

The structure of overlapping subdivergences, which appear in the perturbative expansions of quantum field theory, is analyzed using algebraic lattice theory. It is shown that for specific QFTs the sets of subdivergences of Feynman diagrams form algebraic lattices. This class of QFTs includes the Standard model. In kinematic renormalization schemes, in which *tadpole* diagrams vanish, the lattices are semimodular. This implies that the Hopf algebra of Feynman diagrams is graded by the coradical degree or equivalently that every maximal forest has the same length in the scope of BPHZ renormalization. As an application of this framework a formula for the counter terms in zero-dimensional QFT is given together with some examples of the enumeration of primitive or skeleton diagrams.

1 Introduction

Calculations of observable quantities in quantum field theory rely almost always on perturbation theory. The integrals in the perturbation expansions can be depicted as Feynman diagrams. Usually these integrals require *renormalization* to give meaningful results. The renormalization procedure can be performed using the Hopf algebra of Feynman graphs [7], which organizes the classic BPHZ procedure into an algebraic framework.

The motivation for this paper was to obtain insights on the *coradical filtration* of this Hopf algebra and thereby on the structure of the *subdivergences* of the Feynman diagrams in the perturbation expansion.

The perturbative expansions in QFT are divergent series themselves. This divergence is believed to be dominated by the growth of the number of Feynman diagrams. The coradical filtration describes the hierarchy in which diagrams become important in the large-order regime. Dyson-Schwinger equations exploit this hierarchy to give non-perturbative results [11]. This work also aims to extend the effectiveness of these methods.

Notation and preliminaries on Feynman diagrams and Hopf algebras are covered in section 2 in a form suitable for a combinatorial analysis. Feynman diagrams are defined as a special case of hyper graphs. This definition was used to clarify the role of external legs and isomorphisms of diagrams. Based on this, Kreimer's Hopf algebra of Feynman diagrams is defined.

The starting point in section 3 for the analysis of algebraic lattices in renormalization is the basic fact that *subdivergences* of Feynman diagrams form a partially ordered set or poset ordered by inclusion. In distinguished renormalizable quantum field theories a *join* and a *meet* can be defined on these sets, promoting the posets to *algebraic lattices*. These renormalizable QFTs

*borinsky@physik.hu-berlin.de

will be called *join-meet-renormalizable*. It is shown that a broad class of QFTs including the standard model falls into this category. ϕ^6 -theory in 3-dimensions is examined as an example of a QFT, which is renormalizable, but not join-meet-renormalizable.

The Hopf algebra of decorated posets or lattices is defined, which augments the incidence Hopf algebra [13] by a decoration. Eventually, a Hopf algebra morphism from the Hopf algebra of Feynman diagrams to the Hopf algebra of posets is given.

A further analysis in section 4 demonstrates that in QFTs with only three-or-less-valent vertices, which are thereby join-meet-renormalizable, these lattices are *semimodular*. This implies that the Hopf algebra is bigraded by the loop number of the Feynman diagram and its *coradical degree*. In the language of BPHZ this means that every *complete forest* has the same length. Generally, this structure cannot be found in join-meet-renormalizable theories with also four-valent vertices as QCD or ϕ^4 . An explicit counter example of a non-graded and non-semimodular lattice, which appears in ϕ^4 and Yang-Mills theories, is given. The semimodularity of the subdivergence lattices can be resurrected in these cases by dividing out *tadpole* (also snail or seagull) diagrams. This quotient can always be formed in kinematic renormalization schemes.

This whole framework is used in section 5 to illuminate some results of zero-dimensional QFTs [8, 1] from the perspective of the Hopf algebra of decorated lattices. A closed formula for the counter term calculation is given using the *Moebius function* on the lattices. This formula is applied to the enumeration of primitive diagrams in a variety of cases.

2 Preliminaries

2.1 Combinatorial quantum field theory

In what follows a quantum field theory (QFT) will be characterized by its field content, its interactions, associated ‘weights’ for these interactions and a given dimension of spacetime D . Let Φ denote the set of fields, \mathcal{R}_v the set of allowed interactions, represented as monomials in the fields and $\mathcal{R}_e \subset \mathcal{R}_v$ the set of propagators, a set of distinguished interactions between two fields only. \mathcal{R}_e consists of monomials of degree two and \mathcal{R}_v of monomials of degree two or higher in the fields Φ . Additionally, a map $\omega : \mathcal{R}_e \cup \mathcal{R}_v \rightarrow \mathbb{Z}$ is given associating a weight to each interaction. The requirement $\mathcal{R}_e \subset \mathcal{R}_v$ ensures that there is a two-valent vertex for every allowed edge-type. This is not necessary for the definition of the Hopf algebra of Feynman diagrams, but it results in a simpler formula for contractions among other simplifications. Of course, this does not introduce a restriction to the underlying QFT: A propagator is always associated to the formal inverse of the corresponding two-valent vertex and a two-valent vertex always comes with an additional propagator in a diagram. The two valent vertex of the same type as the propagator can be canceled with the additional propagator.

In physical terms, the interactions correspond to summands in the Lagrangian of the QFT and the weights are the number of derivatives in the respective summand.

Having clarified the important properties of a QFT for a combinatorial treatment, we can proceed to the definition of the central object of perturbative QFTs:

2.2 Feynman diagrams as hypergraphs

Definition 1 (Feynman diagram). A *Feynman diagram* Γ is a tuple (H, E, V, η) . Consisting of

1. a set H of half-edges,

2. a coloring of the half-edges by fields in Φ :

$$\eta : H \rightarrow \Phi. \quad (1)$$

This coloring also induces an additional map, the *residue*,

$$\text{res} : 2^H \rightarrow (\mathbb{N}_0)^\Phi, \quad a \mapsto \prod_{h \in a} \eta(h), \quad (2)$$

of subsets of half-edges, or arbitrary adjacency relations, to monomials in the fields,

3. a set E of edges, adjacency relations between two half-edges with a residue in \mathcal{R}_e ,

$$E \subset \{e \subset H : \text{res}(e) \in \mathcal{R}_e\} \text{ and} \quad (3)$$

4. a set V of vertices or corollas - adjacency relations between any other number of half-edges with a residue in \mathcal{R}_v :

$$V \subset \{v \subset H : \text{res}(v) \in \mathcal{R}_v\}. \quad (4)$$

With the conditions:

5. All the edges and all the vertices are each pairwise disjoint sets: $\forall v_1, v_2 \in V, v_1 \neq v_2 : v_1 \cap v_2 = \emptyset$ as well as $\forall e_1, e_2 \in E, e_1 \neq e_2 : e_1 \cap e_2 = \emptyset$.
6. Every half-edges is in at least one vertex: $\bigcup_{v \in V} v = H$.

Conditions 1-4 form the definition of a *colored hyper graph*. The other conditions 5 and 6 almost restrict these hyper graphs to *multigraphs* with the exception that half-edges that are not in any edge are still allowed. These half-edges will play the role of the external legs of the Feynman diagram. Feynman diagrams will also be called diagrams or graphs in this article. If the reference to the diagram is ambiguous, the sets in the tuple $\Gamma = (H, E, V, \eta)$ will be denoted as $H(\Gamma)$, $E(\Gamma)$ and $V(\Gamma)$.

To clarify the above definition an example is given, in which different depictions of Feynman diagrams are discussed.

Example 1 (Yukawa theory). Let $\Phi = \{\bar{\psi}, \psi, \phi\}$, $\mathcal{R}_v = \{\bar{\psi}\psi, \phi^2, \bar{\psi}\psi\phi\}$ and $\mathcal{R}_e = \{\bar{\psi}\psi, \phi^2\}$. Figure 1 shows different graphical representations for a simple Feynman diagram in this theory.

The usual Feynman diagram representation is given in fig. 1a. The adjacency relations E are represented as edges and the adjacency relations V as vertices. The half-edges are omitted.

Figure 1b shows a hypergraph representation of the diagram. Its half-edges are drawn as little circles. They are colored by the corresponding field. The adjacency relations are shown as big ellipses, enclosing the adjacent half-edges. The adjacency relations, $a \in E \cup V$ can be colored by the different allowed residues, $\text{res}(a)$ in \mathcal{R}_e and \mathcal{R}_v .

In figure 1c, the same diagram is depicted as a bipartite graph. Half-edges are identified with the first class of vertices of the graph and adjacency relations with the second. Half-edges and an adjacency relation are connected if the half-edge appears in the respective adjacency relation. This representation is useful for computational treatments of Feynman diagrams, because colored bipartite graphs are common and well studied in computational graph theory.

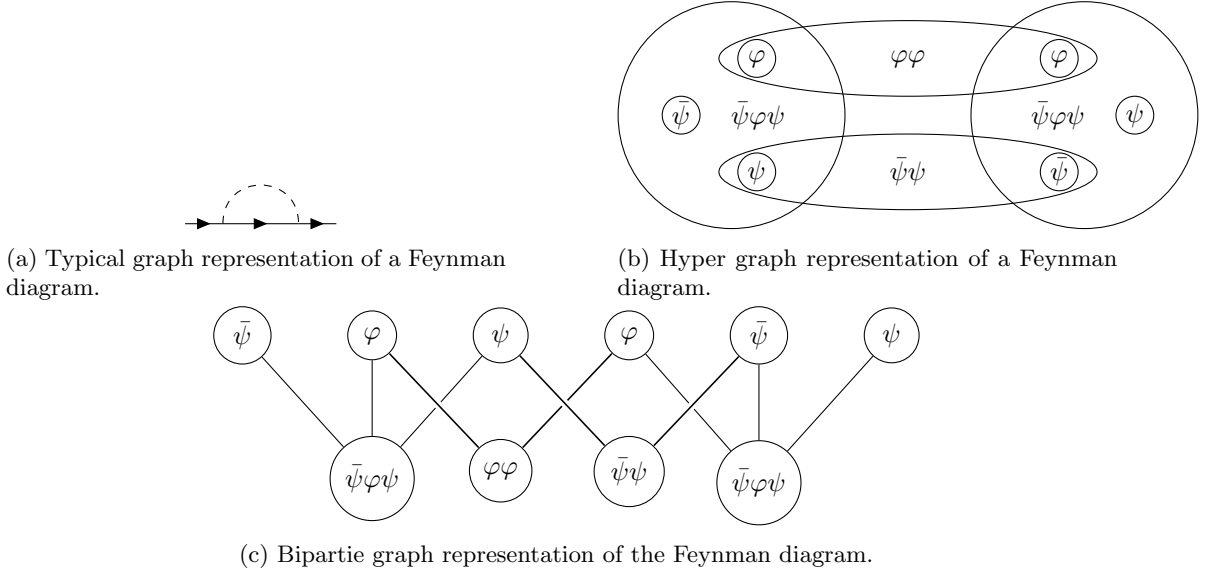


Figure 1: Equivalent diagrammatic representations of Feynman graphs

Isomorphisms of diagrams An *isomorphism* $j : \Gamma \rightarrow \Gamma'$ between one Feynman diagram $\Gamma = (H, V, E, \eta)$ and another $\Gamma' = (H', V', E', \eta')$ is a bijection $j : H \rightarrow H'$, which preserves the adjacency structure $j(E \cup V) = E' \cup V'$ and the coloring $\eta = \eta' \circ j$, where j is extended canonically to $E \cup V \subset 2^H$.

If there is an isomorphism from one diagram Γ to another diagram Γ' , the two diagrams are called isomorphic.

An automorphism is an isomorphism of the diagram onto itself. The group of automorphisms of Γ is denoted as $\text{Aut } \Gamma$. The inverse of the cardinality of the automorphism group, $\frac{1}{|\text{Aut } \Gamma|}$, is the *symmetry factor* of the diagram Γ .

External legs The half-edges which are not included in an edge are called *external legs* or *legs*, $H_{\text{ext}} := H \setminus \bigcup_{e \in E} e$. $H_{\text{int}} = \bigcup_{e \in E} e$ is the set of half-edges which are not external. Feynman diagrams can be classified by the colors of the half-edges in H_{ext} . This classification is called the *residue* of the diagram.

Residues of a diagrams Using the external half-edges, H_{ext} , a monomial $\prod_{h \in H_{\text{ext}}} \eta(h)$ in the fields can be associated to a diagram. This map will also be called *res*:

$$\text{res} : (H, E, V, \eta) \rightarrow (\mathbb{N}_0)^\Phi, \Gamma \mapsto \prod_{h \in H_{\text{ext}}(\Gamma)} \eta(h). \quad (5)$$

The choice of the same name for this map and the coloring of the adjacency relations will be justified by a certain compatibility of *res* for adjacency relations and diagrams in the scope of *renormalizable* QFTs (see definition 2). Furthermore, this should underline the operadic structure of Feynman diagrams.

(Non)-fixed external legs In the above definition the external legs of a diagram do not have a specific labeling as usual. The legs are not ‘fixed’ and permutations of them can result in additional automorphisms. For instance, the diagrams $\times \circ \times \simeq \circ \times \simeq \times \times$ are isomorphic with respect to each other in the present scope. The symmetry factor of $\times \circ \times$ is $\frac{1}{|\text{Aut} \times \circ \times|} = \frac{1}{16}$.

The labeling of the legs was omitted, because it makes the definition of the Hopf algebra of Feynman diagrams more involved without any gain in generality.

Multiplying by the number of possible permutations of the external legs, which is $4! = 24$ in the case of scalar four-leg diagrams, we reobtain the symmetry factors for the leg-fixed case:

$$4! \cdot \frac{1}{16} = \frac{3}{2} = \frac{1}{|\text{Aut}_{\text{lf}} \times \circ \times|} + \frac{1}{|\text{Aut}_{\text{lf}} \circ \times|} + \frac{1}{|\text{Aut}_{\text{lf}} \times \times|}, \quad (6)$$

where Aut_{lf} are the automorphisms which leave the legs fixed. In the same way, we can always replace diagrams with non-fixed legs by a formal sum of diagrams with fixed legs and reobtain the usual notation.

Singular homology For a diagram $\Gamma = (H, E, V, \eta)$, take the free abelian groups with the base field $\mathbb{Z}/2\mathbb{Z}$ with two elements over the sets E and V and let

$$\partial : \mathbb{Z}/2\mathbb{Z}E \rightarrow \mathbb{Z}/2\mathbb{Z}V, \quad e \mapsto \sum_{v \in V} |v \cap e| v, \quad (7)$$

map the edges to the sum over all vertices with each vertex weighted by the number of half-edges which the edge and the vertex share.

The kernel $\ker \partial$ is called the loop space, $\ker \partial =: H_1(\Gamma) \subset \mathbb{Z}/2\mathbb{Z}E$ of the Feynman diagram. Its dimension, $h_1(\Gamma) := \dim H_1(\Gamma)$, is called the *loop number*. This is the first *Betty number* of the topology of the diagram. The cokernel of ∂ is called the space of connected components of the diagram, $\text{coker} \partial =: H_0(\Gamma) \subset \mathbb{Z}/2\mathbb{Z}V$. Its dimension, $h_0(\Gamma) := \dim H_0(\Gamma)$, is the number of connected components. By the Euler characteristic, $h_1(\Gamma) - |E| + |V| - h_0(\Gamma) = 0$.

Subdiagrams A subdiagram of a Feynman diagram $\Gamma = (H, E, V, \eta)$ is a Feynman diagram $\gamma = (H', E', V', \eta')$ such that $E' \cup V' \subset E \cup V$ and $\eta' = \eta|_{H'}$. The relation γ is a subdiagram of Γ is denoted as $\gamma \subset \Gamma$.

The union and the intersection of two subdiagrams $\gamma_1 = (H_1, E_1, V_1, \eta_1)$, $\gamma_2 = (H_2, E_2, V_2, \eta_2) \subset \Gamma = (H, E, V, \eta)$ is defined as

$$\gamma_1 \cup \gamma_2 = (H_1 \cup H_2, E_1 \cup E_2, V_1 \cup V_2, \eta|_{H_1 \cup H_2}) \quad (8)$$

$$\gamma_1 \cap \gamma_2 = (H_1 \cap H_2, E_1 \cap E_2, V_1 \cap V_2, \eta|_{H_1 \cap H_2}). \quad (9)$$

$\gamma_1 \cup \gamma_2$ and $\gamma_1 \cap \gamma_2$ are also subdiagrams of Γ .

Connectedness A diagram Γ is *connected* if it does not have two subdiagrams γ_1 and γ_2 such that $\gamma_1 \cup \gamma_2 = \Gamma$ and $\gamma_1 \cap \gamma_2 = \emptyset$, where \emptyset is the diagram with no half-edges, $|H(\emptyset)| = 0$.

Using the singular homology of Γ , the important notion of 1PI subdiagrams can be introduced:

1PI subdiagrams First, the set of subdiagrams

$$\mathcal{P}_{\text{1PI}}^{L=1}(\Gamma) = \left\{ \gamma \subset \Gamma \text{ such that } \left(\sum_{e \in E(\gamma)} e \right) \in H_1(\Gamma) \text{ and } \gamma \text{ is connected} \right\} \quad (10)$$

is the set of one loop, *1PI subdiagrams* of Γ and

$$\mathcal{P}_{1\text{PI}}(\Gamma) := \left\{ \gamma \subset \Gamma \text{ such that } \gamma = \bigcup_{\gamma' \in I} \gamma' \text{ with } I \subset \mathcal{P}_{1\text{PI}}^{L=1}(\Gamma) \text{ and } \gamma \text{ is connected} \right\} \quad (11)$$

is the general set of 1PI subdiagrams. 1PI stands for *one particle irreducible*. This definition is equivalent to the requirement that the subdiagram is still connected after removing an arbitrary edge. A diagram Γ is called 1PI if $\Gamma \in \mathcal{P}_{1\text{PI}}(\Gamma)$.

Example 2 (1PI subdiagrams of a diagram in ϕ^4 theory). For the diagram $\times \text{---} \bigcirc \text{---} \times$ in ϕ^4 theory ($\Phi = \{\phi\}$, $\mathcal{R}_v = \{\phi^4, \phi^2\}$ and $\mathcal{R}_e = \{\phi^2\}$).

$$\mathcal{P}_{1\text{PI}} \left(\times \text{---} \bigcirc \text{---} \times \right) = \left\{ \begin{array}{l} \times \text{---} \bigcirc \text{---} \times, \times \text{---} \bigcirc \text{---} \times, \times \text{---} \bigcirc \text{---} \times, \times \text{---} \bigcirc \text{---} \times, \times \text{---} \bigcirc \text{---} \times, \times \text{---} \bigcirc \text{---} \times, \\ \times \text{---} \bigcirc \text{---} \times, \times \text{---} \bigcirc \text{---} \times, \times \text{---} \bigcirc \text{---} \times, \times \text{---} \bigcirc \text{---} \times, \times \text{---} \bigcirc \text{---} \times, \times \text{---} \bigcirc \text{---} \times \end{array} \right\},$$

where subdiagrams are drawn with thick lines.

Contractions A 1PI subdiagram $\gamma = (H', E', V', \eta') \in \mathcal{P}_{1\text{PI}}(\Gamma)$ of a diagram $\Gamma = (H, E, V, \eta)$, $\gamma \subset \Gamma$ can be contracted. The result is denoted as Γ/γ . Explicitly:

$$\Gamma/\gamma := ((H \setminus H') \cup H'_{\text{ext}}, E \setminus E', (V \setminus V') \cup \{H'_{\text{ext}}\}, \eta' \big|_{(H \setminus H') \cup H'_{\text{ext}}}). \quad (12)$$

The contraction of a subdiagram involves the removal of all the internal half-edges, all the edges and vertices of the subdiagram. Only the legs of the subdiagram are kept and ‘tied together’ to form a new vertex of the resulting diagram. For a propagator-type subdiagram this means that an additional vertex with the same residue as the contracted graph is added. The new vertex in the produced diagram, $v^* \in \Gamma/\gamma$ has the residue as the contracted subdiagram, $\text{res}(v^*) = \text{res}(\gamma)$. For Γ/γ to be a valid Feynman diagram with vertex-types in \mathcal{R}_v , the residues of contracted subdiagrams need to be restricted in a certain way. This will lead to the notion of superficial divergent subdiagrams. Moreover, this will explain the requirement $\mathcal{R}_e \subset \mathcal{R}_v$.

Superficial degree of divergence Using the map ω , which is provided by the QFT, to assign a weight to every vertex and edge-type, an additional map ω_D can be defined, which assigns a weight to a Feynman diagram. This weight is called *superficial degree of divergence* in the sense of [16]:

$$\omega_D(\Gamma) := \sum_{e \in E} \omega(\text{res}(e)) - \sum_{v \in V} \omega(\text{res}(v)) - Dh_1(\Gamma) \quad (13)$$

Neglecting possible infrared divergences, the value of ω_D coincides with the degree of divergence of the integral associated to the diagram in the perturbation expansion of the underlying QFT in D -dimensions. A 1PI diagram Γ with $\omega_D(\Gamma) \leq 0$ is *superficially divergent* (s.d.) in D dimensions. For notational simplicity, the weight 0 is assigned to the empty diagram, $\omega_D(\emptyset) = 0$, even though it is not divergent.

Definition 2 (Renormalizable Quantum Field Theory). A QFT is *renormalizable in D dimensions* if $\omega_D(\Gamma)$ depends only on the external structure of Γ and the superficial degree of divergence agrees with the weight assigned to the residue of the diagram: $\omega_D(\Gamma) = \omega(\text{res}\Gamma)$. This can be expressed as the commutativity of the diagram:

$$\begin{array}{ccc}
\mathcal{T} & \xrightarrow{\omega_D} & \mathbb{Z} \\
\downarrow \text{res} & \nearrow \omega & \\
\mathbb{N}^\Phi & &
\end{array}$$

where \mathcal{T} is the set of all connected admissible Feynman diagrams of the renormalizable QFT. Specifically, $\omega_D(\Gamma)$ needs to be independent of $h_1(\Gamma)$. Working with a renormalizable QFT, we need to keep track of *subdivergences* or *superficially divergent subdiagrams* appearing in the integrals of the perturbation expansion. The tools needed are the set of 1PI subdiagrams and the superficial degree of divergence. The compatibility of the vertex and edge-weights and the superficial degree of divergence of the diagrams is exactly what is necessary to contract these subdivergences without leaving the space of allowed Feynman diagrams.

Superficially divergent subdiagrams The set of *superficially divergent subdiagrams* or s.d. subdiagrams,

$$\mathcal{P}_D^{\text{s.d.}}(\Gamma) := \left\{ \gamma \subset \Gamma \text{ such that } \gamma = \prod_i \gamma_i, \gamma_i \in \mathcal{P}_{\text{1PI}}(\Gamma) \text{ and } \omega_D(\gamma_i) \leq 0 \right\}, \quad (14)$$

of subdiagrams, whose connected components are s.d. 1PI diagrams, is the object of main interest for the combinatorics of renormalization. The renormalizability of the QFT guarantees that for every $\gamma \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$ the diagram resulting from the contraction Γ/γ is still a valid Feynman diagram of the underlying QFT.

Example 3 (Superficially divergent subdiagrams of a diagram in ϕ^4 theory). Consider the same diagram as in example 2 in ϕ^4 theory with the weights $\omega(\phi^2) = \omega(-) = 2$ and $\omega(\phi^4) = \omega(\times) = 0$. The superficially divergent subdiagrams for $D = 4$ are

$$\mathcal{P}_4^{\text{s.d.}} \left(\text{diagram} \right) = \left\{ \text{diagram}_1, \text{diagram}_2, \text{diagram}_3, \text{diagram}_4, \text{diagram}_5 \right\}.$$

2.3 Hopf algebra structure of Feynman diagrams

The basis for the analysis of the lattice structure in QFTs is Kreimer's Hopf algebra of Feynman diagrams. It captures the BPHZ renormalization procedure which is necessary to obtain finite amplitudes from perturbative calculations in an algebraic framework [6]. In this section, the basic definitions of the Hopf algebra of Feynman diagrams will be repeated. For a more detailed exposition consult [12] for mathematical details of Hopf algebras or [3] for more computational aspects.

Definition 3. Let $\mathcal{H}_D^{\text{fg}}$ be the \mathbb{Q} -algebra generated by all mutually non-isomorphic Feynman diagrams, whose connected components are superficially divergent 1PI diagrams of a certain QFT in D -dimensions. The multiplication of generators is given by the disjoint union: $m : \mathcal{H}_D^{\text{fg}} \otimes \mathcal{H}_D^{\text{fg}} \rightarrow \mathcal{H}_D^{\text{fg}}, \gamma_1 \otimes \gamma_2 \mapsto \gamma_1 \cup \gamma_2$. It is extended linearly to all elements in the vector space $\mathcal{H}_D^{\text{fg}}$. $\mathcal{H}_D^{\text{fg}}$ has a unit $u : \mathbb{Q} \mapsto \mathbb{I} \subset \mathcal{H}_D^{\text{fg}}$, where \mathbb{I} is associated to the empty diagram and a counit $\epsilon : \mathcal{H}_D^{\text{fg}} \rightarrow \mathbb{Q}$, which vanishes on every generator of $\mathcal{H}_D^{\text{fg}}$ except \mathbb{I} : $\epsilon \circ \mathbb{I} := 1$. The coproduct on the generators is defined as follows:

$$\Delta_D \Gamma := \sum_{\gamma \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)} \gamma \otimes \Gamma/\gamma \quad : \quad \mathcal{H}_D^{\text{fg}} \rightarrow \mathcal{H}_D^{\text{fg}} \otimes \mathcal{H}_D^{\text{fg}}, \quad (15)$$

where the complete contraction Γ/Γ is set to \mathbb{I} and the vacuous contraction Γ/\emptyset to Γ . The notion of superficial degree of divergence, ω_D , hidden in $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ (eq. (14)) is the only input to the Hopf algebra structure which depends on the dimension D of spacetime.

Example 4 (Coproduct of a diagram in ϕ^4 -theory). Take the same diagram of ϕ^4 -theory as in examples 2 and 3. The coproduct is calculated using the set $\mathcal{P}_4^{\text{s.d.}}(\Gamma)$ and the definition of the contraction in eq. (12):

$$\begin{aligned} \Delta_4 \times \text{diagram} &= \sum_{\gamma \in \{\text{diagram}_1, \text{diagram}_2, \text{diagram}_3, \text{diagram}_4, \text{diagram}_5\}} \gamma \otimes \text{diagram} / \gamma = \\ &\mathbb{I} \otimes \text{diagram} + \text{diagram} \otimes \mathbb{I} + \\ &+ \text{diagram}_1 \otimes \text{diagram}_2 + \text{diagram}_2 \otimes \text{diagram}_1 + \text{diagram}_3 \otimes \text{diagram}_4 + \text{diagram}_4 \otimes \text{diagram}_3 = \\ &\mathbb{I} \otimes \text{diagram} + \text{diagram} \otimes \mathbb{I} + 2 \times \text{diagram}_1 \otimes \text{diagram}_2 + \text{diagram}_3 \otimes \text{diagram}_4 \end{aligned}$$

The last equality holds because diagram_1 and diagram_2 are the same diagrams up to a permutation of the legs.

Extended linearly to all elements of $\mathcal{H}_D^{\text{fg}}$, the coproduct $\Delta_D : \mathcal{H}_D^{\text{fg}} \rightarrow \mathcal{H}_D^{\text{fg}} \otimes \mathcal{H}_D^{\text{fg}}$ is an algebra morphism. This implies $\Delta_D \mathbb{I} = \mathbb{I} \otimes \mathbb{I}$ and $\Delta_D \circ m = (m \otimes m) \circ \tau_{23} \circ (\Delta_D \otimes \Delta_D)$, where τ_{23} switches the second and the third entry of the tensor product. The coproduct is coassociative: $(\Delta_D \otimes \text{id}) \circ \Delta_D = (\text{id} \otimes \Delta_D) \circ \Delta_D$. $\mathcal{H}_D^{\text{fg}}$ is graded by the loop number, $h_1(\Gamma)$, of the diagrams,

$$\mathcal{H}_D^{\text{fg}} = \bigoplus_{L \geq 0} \mathcal{H}_D^{\text{fg}(L)} \quad \text{and} \quad (16)$$

$$m : \mathcal{H}_D^{\text{fg}(L_1)} \otimes \mathcal{H}_D^{\text{fg}(L_2)} \rightarrow \mathcal{H}_D^{\text{fg}(L_1+L_2)} \quad (17)$$

$$\Delta_D : \mathcal{H}_D^{\text{fg}(L)} \rightarrow \bigoplus_{\substack{L_1, L_2 \geq 0 \\ L_1+L_2=L}} \mathcal{H}_D^{\text{fg}(L_1)} \otimes \mathcal{H}_D^{\text{fg}(L_2)}, \quad (18)$$

where $\mathcal{H}_D^{\text{fg}(L)} \subset \mathcal{H}_D^{\text{fg}}$ is the subspace of $\mathcal{H}_D^{\text{fg}}$ which is generated by diagrams Γ with $h_1(\Gamma) = L$. Because $\mathcal{H}_D^{\text{fg}}$ is coassociative, it makes sense to define iterations of Δ_D :

$$\Delta_D^0 := \epsilon, \quad \Delta_D^1 := \text{id}, \quad \Delta_D^n := (\Delta_D \otimes \text{id}^{\otimes n-2}) \circ \Delta_D^{n-1} \quad \text{for } n \geq 2, \quad (19)$$

where $\Delta_D^k : \mathcal{H}_D^{\text{fg}} \rightarrow \mathcal{H}_D^{\text{fg}^{\otimes k}}$. $\mathcal{H}_D^{\text{fg}^{\otimes 0}} \simeq \mathbb{Q}$, because the unit \mathbb{I} is the only generator with 0 loops. For this reason, $\mathcal{H}_D^{\text{fg}}$ is a connected Hopf algebra. The reduced coproduct is defined as $\tilde{\Delta}_D := P^{\otimes 2} \circ \Delta_D$, where $P := \text{id} - u \circ \epsilon$ projects into the augmentation ideal, $P : \mathcal{H}_D^{\text{fg}} \rightarrow \bigoplus_{L \geq 1} \mathcal{H}_D^{\text{fg}(L)}$.

Example 5 (Reduced coproduct of a non-primitive diagram in ϕ^4 -theory).

$$\tilde{\Delta}_4 \times \text{diagram} = 2 \times \text{diagram}_1 \otimes \text{diagram}_2 + \text{diagram}_3 \otimes \text{diagram}_4 \quad (20)$$

The kernel of the reduced coproduct, is the space of *primitive* elements of the Hopf algebra, $\text{Prim} \mathcal{H}_D^{\text{fg}} := \ker \tilde{\Delta}_D$. Primitive 1PI diagrams Γ with $\Gamma \in \ker \tilde{\Delta}_D$ are exactly those diagrams, which do not contain any subdivergences. They are also called *skeleton diagrams*. More general,



Figure 2: Characteristic subdiagram of every tadpole diagram

we can define $\tilde{\Delta}_D^n = P^{\otimes n} \circ \Delta_D^n$. These morphisms give rise to an increasing filtration of $\mathcal{H}_D^{\text{fg}}$, the *coradical filtration*:

$${}^{(n)}\mathcal{H}_D^{\text{fg}} := \ker \tilde{\Delta}_D^{n+1} \quad \forall n \geq 0 \quad (21)$$

$$\mathbb{Q} \simeq {}^{(0)}\mathcal{H}_D^{\text{fg}} \subset {}^{(1)}\mathcal{H}_D^{\text{fg}} \subset \dots \subset {}^{(n)}\mathcal{H}_D^{\text{fg}} \subset \dots \subset \mathcal{H}_D^{\text{fg}}. \quad (22)$$

$\mathcal{H}_D^{\text{fg}}$ is equipped with an *antipode* $S_D : \mathcal{H}_D^{\text{fg}} \rightarrow \mathcal{H}_D^{\text{fg}}$, implicitly defined by the identity, $m \circ (\text{id} \otimes S_D) \circ \Delta_D = \text{u} \circ \epsilon$. Because $\mathcal{H}_D^{\text{fg}}$ is connected, there is always a unique solution for S_D , which can be calculated recursively.

Group of characters The characters of $\mathcal{H}_D^{\text{fg}}$ are linear algebra morphisms from $\mathcal{H}_D^{\text{fg}}$ to some commutative, unital algebra \mathcal{A} . The set of all these morphisms forms a group and is denoted as $G_{\mathcal{A}}^{\mathcal{H}_D^{\text{fg}}}$. The product on this group, called the *convolution product*, is given by $\phi * \psi = m_{\mathcal{A}} \circ (\phi \otimes \psi) \circ \Delta_D$ for $\phi, \psi \in G_{\mathcal{A}}^{\mathcal{H}_D^{\text{fg}}}$ and the unit of the group is the morphism $\text{u}_{\mathcal{A}} \circ \epsilon_{\mathcal{H}_D^{\text{fg}}}$.

Hopf ideals A *Hopf ideal* I of a Hopf algebra \mathcal{H} is an ideal of the algebra and a coideal of the coalgebra. Explicitly, it is a subspace of a Hopf algebra, $I \subset \mathcal{H}$, such that

$$m(I \otimes \mathcal{H}) \subset I \quad (23)$$

$$\Delta I \subset I \otimes \mathcal{H} + \mathcal{H} \otimes I. \quad (24)$$

For every ideal I , the quotient Hopf algebra \mathcal{H}/I can be defined, where two elements $x, y \in \mathcal{H}$ are equivalent if their difference is in the ideal $x - y \in I$. The quotient \mathcal{H}/I is also a Hopf algebra.

Hopf algebra morphisms Algebra morphisms which also respect the coalgebraic structure of the Hopf algebra are called *Hopf algebra morphisms*. They are linear maps from a Hopf algebra \mathcal{H}_1 to another Hopf algebra \mathcal{H}_2 . Explicitly, this means that $\psi : \mathcal{H}_1 \rightarrow \mathcal{H}_2$ is a Hopf algebra morphism iff $m_{\mathcal{H}_2} \circ (\psi \otimes \psi) = \psi \circ m_{\mathcal{H}_1}$ and $\Delta_{\mathcal{H}_2} \circ \psi = (\psi \otimes \psi) \circ \Delta_{\mathcal{H}_1}$. A Hopf algebra morphism ψ always gives rise to a Hopf ideal in \mathcal{H}_1 . This ideal is the kernel of the morphism denoted as $\ker \psi$. The quotient Hopf algebra $\mathcal{H}_1 / \ker \psi$ is isomorphic to \mathcal{H}_2 .

An useful example of a Hopf algebra morphism is the projection of $\mathcal{H}_D^{\text{fg}}$ to the Hopf algebra of Feynman diagrams without ‘tadpoles’ (also snails or seagulls). Tadpoles are diagrams which can be split into two connected components by removing a single vertex such that one component does not contain any external leg. A tadpole diagram always has a subdiagram of a topology as depicted in fig. 2. The Hopf algebra of Feynman diagrams without tadpoles is denoted as $\tilde{\mathcal{H}}_D^{\text{fg}}$.

Definition 4. We define $\tilde{\mathcal{H}}_D^{\text{fg}}$ as $\mathcal{H}_D^{\text{fg}}$ with the difference that no tadpole diagrams are allowed as generators and replace $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ in the formula for the coproduct, eq. (15), with

$$\tilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma) := \{\gamma \in \mathcal{P}_D^{\text{s.d.}}(\Gamma) : \text{such that } \Gamma/\gamma \text{ is no tadpole diagram}\}. \quad (25)$$

Only the s.d. subdiagrams which do not result in a tadpole diagram upon contraction are elements of $\widetilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma)$.

A Hopf algebra morphism from $\mathcal{H}_D^{\text{fg}}$ to $\widetilde{\mathcal{H}}_D^{\text{fg}}$ is easy to set up:

$$\psi \quad : \quad \mathcal{H}_D^{\text{fg}} \quad \rightarrow \quad \widetilde{\mathcal{H}}_D^{\text{fg}}, \quad (26)$$

$$\Gamma \quad \mapsto \quad \begin{cases} 0 & \text{if } \Gamma \text{ is a tadpole diagram.} \\ \Gamma & \text{else.} \end{cases} \quad (27)$$

This map fulfills the requirements for a Hopf algebra morphism. The associated ideal $\ker \psi \subset \mathcal{H}_D^{\text{fg}}$ is the subspace of $\mathcal{H}_D^{\text{fg}}$ spanned by all tadpole diagrams. This ideal and the map ψ are very useful, because the elements in $\ker \psi$ evaluate to zero after renormalization in kinematic subtraction schemes [5] and in minimal subtraction schemes for the massless case.

3 Algebraic lattice structure of subdivergences

To highlight the lattice structure of renormalizable QFTs, it is necessary to have more information on the structure in which subdivergences can appear. The classic notion of overlapping divergence, which addresses the case of two or more subdiagrams which are neither contained in each other nor disjoint, is central for this analysis.

3.1 Overlapping divergences

Definition 5. Two diagrams $\gamma_1, \gamma_2 \in \mathcal{P}_{1\text{PI}}(\Gamma)$ are called overlapping if $\gamma_1 \not\subset \gamma_2$, $\gamma_2 \not\subset \gamma_1$ and $\gamma_1 \cap \gamma_2 \neq \emptyset$.

To proof the important properties of overlapping divergences exploited in the following sections, we need the following two lemmas:

Lemma 1. For two arbitrary subdiagrams $\gamma_1, \gamma_2 \subset \Gamma$:

$$H_{\text{ext}}(\gamma_1 \cap \gamma_2) = (H_{\text{ext}}(\gamma_1) \cup H_{\text{ext}}(\gamma_2)) \cap H(\gamma_1) \cap H(\gamma_2) \quad (28)$$

$$H_{\text{ext}}(\gamma_1 \cup \gamma_2) = (H_{\text{ext}}(\gamma_1) \cup H_{\text{ext}}(\gamma_2)) \setminus H_{\text{int}}(\gamma_1) \setminus H_{\text{int}}(\gamma_2) \quad (29)$$

$$|H_{\text{ext}}(\gamma_1)| + |H_{\text{ext}}(\gamma_2)| = |H_{\text{ext}}(\gamma_1 \cup \gamma_2)| + |H_{\text{ext}}(\gamma_1 \cap \gamma_2)|. \quad (30)$$

Proof. We abbreviate $H(\gamma_1) =: H^1$, $H_{\text{int}}(\gamma_2) =: H_i^2$, $H_{\text{ext}}(\gamma_1 \cup \gamma_2) =: H_e^{1 \cup 2}$, $H_{\text{ext}}(\gamma_1 \cap \gamma_2) =: H_e^{1 \cap 2}$ and so on. The first and the second identity follow from the definition of $H_{\text{ext}}(\gamma)$ and the condition that edges of a diagram Γ are either disjoint or equal:

$$H_e^{1 \cap 2} = (H^1 \cap H^2) \setminus \left(\bigcup_{e \in E(\gamma_1) \cap E(\gamma_2)} e \right) \quad (31)$$

using condition 5 of definition 1,

$$= (H^1 \cap H^2) \setminus (H_i^1 \cap H_i^2) = (H_e^1 \cap H^2) \cup (H_e^2 \cap H^1) = (H_e^1 \cup H_e^2) \cap H^1 \cap H^2. \quad (32)$$

Analogously, the second identity follows. The third identity follows from the other two by applying inclusion-exclusion on the right hand side of eq. (28) and eq. (29), $|H_e^{1 \cap 2}| = |H_e^1 \cap$

$H^2| + |H_e^2 \cap H^1| - |H_e^1 \cap H_e^2|$, $|H_e^{1 \cup 2}| = |H_e^1 \setminus H_e^2| + |H_e^2 \setminus H_e^1| - |H_e^1 \cap H_e^2|$, adding both and using inclusion-exclusion again:

$$\begin{aligned} |H_e^{1 \cap 2}| + |H_e^{1 \cup 2}| &= |(H_e^1 \cap H^2) \cup (H_e^1 \setminus H_e^2)| + |(H_e^1 \cap H^2) \cap (H_e^1 \setminus H_e^2)| \\ &\quad + |(H_e^2 \cap H^1) \cup (H_e^2 \setminus H_e^1)| + |(H_e^2 \cap H^1) \cap (H_e^2 \setminus H_e^1)| \\ &\quad - 2|H_e^1 \cap H_e^2| = |H_e^1| + |H_e^2| \end{aligned} \quad (33)$$

□

Lemma 2. For every proper subdiagram γ of a connected diagram Γ :

$$|H_{\text{ext}}(\gamma) \cap H_{\text{int}}(\Gamma)| \geq 1 \quad (34)$$

Proof. Suppose $|H_{\text{ext}}(\gamma) \cap H_{\text{int}}(\Gamma)| = 0$, that means there are no external legs of γ , which are internal in Γ . As a consequence, all half-edges of γ can be removed from Γ without breaking an edge. That means, $\Gamma \setminus \gamma = (H(\Gamma) \setminus H(\gamma), E(\Gamma) \setminus E(\gamma), V(\Gamma) \setminus V(\gamma), \eta|_{H(\Gamma) \setminus H(\gamma)})$ is also a Feynman diagram. But $(\Gamma \setminus \gamma) \cup \gamma = \Gamma$ and $(\Gamma \setminus \gamma) \cap \gamma = \emptyset$, thus Γ must be disconnected. □

Corollary 1. For every proper subdiagram γ of a 1PI diagram Γ :

$$|H_{\text{ext}}(\gamma) \cap H_{\text{int}}(\Gamma)| \geq 2 \quad (35)$$

Proof. Suppose $|H_{\text{ext}}(\gamma) \cap H_{\text{int}}(\Gamma)| = 1$ and let $h \in H_{\text{ext}}(\gamma) \cap H_{\text{int}}(\Gamma)$. As h belongs to some edge $e \in E$ of $\Gamma = (H, E, V, \eta)$ such that $h \in e$, this edge can be removed from Γ . Let $\Gamma' = (H, E \setminus e, V, \eta)$ which is still a connected diagram, because Γ is 1PI. Accordingly, $|H_{\text{ext}}(\gamma) \cap H_{\text{int}}(\Gamma')| = 0$, which contradicts the result of lemma 2. □

Proposition 1. If γ_1 and γ_2 are overlapping subdiagrams in $\mathcal{P}_{1\text{PI}}(\Gamma)$, and $\mu \subset \gamma_1 \cap \gamma_2$ some connected component of $\gamma_1 \cap \gamma_2$, then $|H_{\text{ext}}(\mu)| \geq 4$.

Proof. γ_1 and γ_2 are overlapping and μ is a proper subdiagram of both. In this case *corollary1* applies: $|H_{\text{ext}}(\mu) \cap H_{\text{int}}(\gamma_1)| \geq 2$ and $|H_{\text{ext}}(\mu) \cap H_{\text{int}}(\gamma_2)| \geq 2$. It remains to be proven that the sets $H_{\text{ext}}(\mu) \cap H_{\text{int}}(\gamma_1)$ and $H_{\text{ext}}(\mu) \cap H_{\text{int}}(\gamma_2)$ are mutually disjoint. μ is a connected component of $\gamma_1 \cap \gamma_2$, consequently $H_{\text{ext}}(\mu) \subset H_{\text{ext}}(\gamma_1 \cap \gamma_2)$. Lemma 1 implies $H_{\text{ext}}(\gamma_1 \cap \gamma_2) \subset H_{\text{ext}}(\gamma_1) \cup H_{\text{ext}}(\gamma_2) \subset H_{\text{int}}(\gamma_1) \cap H_{\text{int}}(\gamma_2)$. Hence, $H_{\text{ext}}(\mu) \cap H_{\text{int}}(\gamma_1) \cap H_{\text{int}}(\gamma_2) = \emptyset$. □

3.2 Posets and algebraic lattices

The set of subdivergences of a Feynman diagram is obviously partially ordered by inclusion. These posets are quite constrained for some renormalizable QFTs: They additionally carry a lattice structure. In this section, the necessary definitions of poset and lattice theory will be repeated followed by the definition of *join-meet-renormalizability* and the introduction of the Hopf algebra of decorated posets. The definitions will be illustrated by the application to the set of subdivergences of a Feynman diagram. For a more detailed exposition of poset and lattice theory we refer to [14].



Definition 6 (Poset). A *partially ordered set* or *poset* is a finite set P endowed with a partial order \leq . An interval $[x, y]$ is a subset $\{z \in P : x \leq z \leq y\} \subset P$. If $[x, y] = \{x, y\}$, x covers y and y is covered by x .

Hasse diagram A Hasse diagram of a poset P is the graph with the elements of P as vertices and the cover relations as edges. Larger elements are always drawn above smaller elements.

Example 6. The set of superficially divergent subdiagrams $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ of a Feynman diagram Γ is a poset ordered by inclusion: $\gamma_1 \leq \gamma_2 \Leftrightarrow \gamma_1 \subset \gamma_2$ for all $\gamma_1, \gamma_2 \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$.

The statement that a subdiagram γ_1 covers γ_2 in $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ is equivalent to the statement that γ_1/γ_2 is primitive. The elements that are covered by the full diagram $\Gamma \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$ are called *maximal forests*. Whereas, a maximal chain $\emptyset \subset \gamma_1 \subset \dots \subset \gamma_n \subset \Gamma$, where each element is covered by the next, is a *complete forest* of Γ .

The Hasse diagram of a s.d. diagram Γ can be constructed by the following procedure: Draw the diagram and find all the maximal forests $\gamma_i \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$ such that Γ/γ_i is primitive. Draw the diagrams γ_i under Γ and draw lines from Γ to the γ_i . Subsequently, determine all the maximal forests μ_i of the γ_i and draw them under the γ_i . Draw a line from γ_i to μ_i if $\mu_i \subset \gamma_i$. Repeat this until only primitive diagrams are left. Then draw lines from the primitive subdiagrams to an additional \emptyset -diagram underneath them. In the end, replace diagrams by vertices.

Example 7. For instance, the set of superficially divergent subdiagrams for $D = 4$ of the diagram,  can be represented as the Hasse diagram , where the vertices represent the subdiagrams in the set given in example 3.

Definition 7 (Lattice). A lattice is a poset L for which an unique least upper bound (*join*) and an unique greatest lower bound (*meet*) exists for any combination of two elements in L . The join of two elements $x, y \in L$ is denoted as $x \vee y$ and the meet as $x \wedge y$. Every lattice has a unique greatest element denoted as $\hat{1}$ and a unique smallest element $\hat{0}$. Every interval of a lattice is also a lattice.

In many QFTs $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ is a lattice for every s.d. diagram Γ :

Definition 8 (Join-meet-renormalizable quantum field theory). A renormalizable QFT is called join-meet-renormalizable if $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$, ordered by inclusion, is a lattice for every s.d. Feynman diagram Γ .

Theorem 1. A renormalizable QFT is join-meet-renormalizable if $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ is closed under taking unions: $\gamma_1, \gamma_2 \in \mathcal{P}_D^{\text{s.d.}}(\Gamma) \Rightarrow \gamma_1 \cup \gamma_2 \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$ for all s.d. diagrams Γ .

Proof. $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ is ordered by inclusion $\gamma_1 \leq \gamma_2 \Leftrightarrow \gamma_1 \subset \gamma_2$. The join is given by taking the union of diagrams: $\gamma_1 \vee \gamma_2 := \gamma_1 \cup \gamma_2$. $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ has a unique greatest element $\hat{1} := \Gamma$ and a unique smallest element $\hat{0} := \emptyset$. Therefore $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ is a lattice [14, Prop. 3.3.1]. The unique meet is given by the formula, $\gamma_1 \wedge \gamma_2 := \bigcup_{\mu \leq \gamma_1 \text{ and } \mu \leq \gamma_2} \mu$. \square

A broad class of renormalizable QFTs is join-meet-renormalizable including the standard model:

Theorem 2. If all diagrams with four or more legs in a renormalizable QFT are superficially logarithmic divergent or superficially convergent, then the underlying QFT is join-meet-renormalizable.

Proof. From $\gamma_1, \gamma_2 \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$ immediately follows $\gamma_1 \cup \gamma_2 \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$ if γ_1 and γ_2 are disjoint or contained in each other. The statement only needs to be validated if γ_1 and γ_2 are overlapping. In this case, we have $\gamma_1 \cup \gamma_2 \in \mathcal{P}_{\text{IFI}}(\Gamma)$.



(a) Example of a diagram where $\mathcal{P}_3^{\text{s.d.}}(\Gamma)$ is not a lattice. (b) The corresponding non-lattice poset. Trivial vertex multiplicities were omitted.

Figure 3: Counter-example for a renormalizable but not join-meet-renormalizable QFT: ϕ^6 -theory in 3 dimensions.

From definition 2.2 of ω_D and inclusion-exclusion:

$$\omega_D(\gamma_1 \cup \gamma_2) = \omega_D(\gamma_1) + \omega_D(\gamma_2) - \omega_D(\gamma_1 \cap \gamma_2) \quad (36)$$

$\gamma_1 \cap \gamma_2$ has four or more external legs as was proven in proposition 1. Consequently, $\omega_D(\gamma_1 \cap \gamma_2) \geq 0$ and if $\omega_D(\gamma_1) \leq 0$ and $\omega_D(\gamma_2) \leq 0$, then $\omega_D(\gamma_1 \cup \gamma_2) \leq 0$. For this reason $\gamma_1 \cup \gamma_2$ is superficially divergent, $\gamma_1 \cup \gamma_2 \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$ and $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ is closed under taking unions. \square

Corollary 2. *All renormalizable QFTs with only four-or-less-valent vertices are join-meet-renormalizable.*

Proof. If there are only four-or-less-valent vertices, the diagrams with $|H_{\text{ext}}(\Gamma)| > 4$ must be superficially convergent. This also implies that $\omega_D(\Gamma) \geq 0$ for $|H_{\text{ext}}(\Gamma)| \geq 4$. \square

In general, renormalizable QFTs are not join-meet-renormalizable. Fig. 3 shows an example of a s.d. diagram Γ , where $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ is not a lattice. The diagram is depicted in fig. 3a and the corresponding poset in fig. 3b. The diagram appears in ϕ^6 -theory, which is therefore renormalizable, but not join-meet-renormalizable, in 3-dimensions.

To proceed to the Hopf algebra of decorated posets some additional notation of poset and lattice theory must be introduced:

Order preserving maps A map $\sigma : P \rightarrow \mathbb{N}_0$ on a poset to the non-negative numbers is called strictly order preserving if $x < y$ implies $\sigma(x) < \sigma(y)$ for all $x, y \in P$.

Cartesian product of posets From two posets P_1 and P_2 a new poset $P_1 \times P_2 = \{(s, t) : s \in P_1 \text{ and } t \in P_2\}$, the Cartesian product, with the order relation, $(s, t) \leq (s', t')$ iff $s \leq s'$ and $t \leq t'$, is obtained.

The Cartesian product is commutative and if P_1 and P_2 are lattices $P_1 \times P_2$ is also a lattice [14]. It is compatible with the notion of intervals:

$$P_1 \times P_2 \supset [(s, t), (s', t')] = \{(x, y) \in P_1 \times P_2 : s \leq x \leq s' \wedge t \leq y \leq t'\} = [s, s'] \times [t, t']. \quad (37)$$

Isomorphisms of posets An isomorphism between two posets P_1 and P_2 is a bijection $j : P_1 \rightarrow P_2$, which preserves the order relation: $j(x) \leq j(y) \Leftrightarrow x \leq y$.

3.3 The Hopf algebra of decorated posets

Using the preceding notions a new Hopf algebra structure on posets, suitable for the description of the subdivergences, can be defined. This structure is essentially the one of an incidence Hopf algebra [13] augmented by a strictly order preserving map as a decoration.

Definition 9 (Hopf algebra of decorated posets). Let \mathcal{D} be the set of tuples (P, ν) , where P is a finite poset with a unique lower bound $\hat{0}$ and a unique upper bound $\hat{1}$ and a strictly order preserving map $\nu : P \rightarrow \mathbb{N}_0$ with $\nu(\hat{0}) = 0$. One can think of \mathcal{D} as the set of bounded posets augmented by a strictly order preserving decoration. An equivalence relation is set up on \mathcal{D} by relating $(P_1, \nu_1) \sim (P_2, \nu_2)$ if there is an isomorphism $j : P_1 \rightarrow P_2$, which respects the decoration ν : $\nu_1 = \nu_2 \circ j$.

Let \mathcal{H}^P be the \mathbb{Q} -algebra generated by all the elements in the quotient \mathcal{P}/\sim with the commutative multiplication:

$$m_{\mathcal{H}^P} : \quad \mathcal{H}^P \otimes \mathcal{H}^P \quad \rightarrow \quad \mathcal{H}^P, \quad (38)$$

$$(P_1, \nu_1) \otimes (P_2, \nu_2) \quad \mapsto \quad (P_1 \times P_2, \nu_1 + \nu_2), \quad (39)$$

which takes the Cartesian product of the two posets and adds the decorations ν . The sum of the two functions ν_1 and ν_2 is to be interpreted in the sense: $(\nu_1 + \nu_2)(x, y) = \nu_1(x) + \nu_2(y)$. The singleton poset $P = \{\hat{0}\}$ with $\hat{0} = \hat{1}$ and the trivial decoration $\nu(\hat{0}) = 0$ serves as a multiplicative unit: $u(1) = \mathbb{1}_{\mathcal{H}^P} := (\{\hat{0}\}, \hat{0} \mapsto 0)$.

Equipped with the coproduct,

$$\Delta_{\mathcal{H}^P} : \quad \mathcal{H}^P \quad \rightarrow \quad \mathcal{H}^P \otimes \mathcal{H}^P, \quad (40)$$

$$(P, \nu) \quad \mapsto \quad \sum_{x \in P} ([\hat{0}, x], \nu) \otimes ([x, \hat{1}], \nu - \nu(x)), \quad (41)$$

where $(\nu - \nu(x))(y) = \nu(y) - \nu(x)$ and the counit ϵ which vanishes on every generator except $\mathbb{1}_{\mathcal{H}^P}$, the algebra \mathcal{H}^P becomes a counital coalgebra.

Proposition 2. \mathcal{H}^P is a bialgebra.

Proof. The compatibility of the multiplication with the coproduct needs to be proven. Let $(P_1, \nu_1), (P_2, \nu_2) \in \mathcal{P}$.

$$\begin{aligned} \Delta_{\mathcal{H}^P} \circ m_{\mathcal{H}^P}((P_1, \nu_1) \otimes (P_2, \nu_2)) &= \Delta_{\mathcal{H}^P}(P_1 \times P_2, \nu_1 + \nu_2) = \\ &= \sum_{x \in P_1 \times P_2} ([\hat{0}, x], \nu_1 + \nu_2) \otimes ([x, \hat{1}], \nu_1 + \nu_2 - \nu_1(x) - \nu_2(x)) = \\ &= \sum_{y \in P_1} \sum_{z \in P_2} ([\hat{0}_{P_1}, y] \times [\hat{0}_{P_2}, z], \nu_1 + \nu_2) \otimes ([y, \hat{1}_{P_1}] \times [z, \hat{1}_{P_2}], \nu_1 + \nu_2 - \nu_1(x) - \nu_2(x)) = \\ &= (m_{\mathcal{H}^P} \otimes m_{\mathcal{H}^P}) \circ \sum_{y \in P_1} \sum_{z \in P_2} ([[\hat{0}_{P_1}, y], \nu_1) \otimes ([[\hat{0}_{P_2}, z], \nu_2) \\ &\quad \otimes ([y, \hat{1}_{P_1}], \nu_1 - \nu_1(x)) \otimes ([z, \hat{1}_{P_2}], \nu_2 - \nu_2(x))] = \\ &= (m_{\mathcal{H}^P} \otimes m_{\mathcal{H}^P}) \circ \tau_{2,3} \circ (\Delta_{\mathcal{H}^P} \otimes \Delta_{\mathcal{H}^P})((P_1, \nu_1) \otimes (P_2, \nu_2)), \end{aligned} \quad (42)$$

where $\tau_{2,3}$ switches the second and the third factor of the tensor product. \square

Corollary 3. \mathcal{H}^P is a connected Hopf algebra.

Proof. \mathcal{H}^P is graded by the value of $\nu(\hat{1})$. There is only one element of degree 0 because ν must be strictly order preserving. It follows that \mathcal{H}^P is a graded, connected bialgebra and therefore a Hopf algebra [12]. \square

3.4 A Hopf algebra morphism from Feynman diagrams to lattices

Theorem 3. *Let $\nu(\gamma) = h_1(\gamma)$. The map,*

$$\chi_D : \begin{array}{ccc} \mathcal{H}_D^{\text{fg}} & \rightarrow & \mathcal{H}^{\text{P}}, \\ \Gamma & \mapsto & (\mathcal{P}_D^{\text{s.d.}}(\Gamma), \nu), \end{array} \quad (43)$$

$$\Gamma \mapsto (\mathcal{P}_D^{\text{s.d.}}(\Gamma), \nu), \quad (44)$$

which assigns to every diagram, its poset of s.d. subdiagrams decorated by the loop number of the subdiagram, is a Hopf algebra morphism.

Proof. First, it needs to be shown that χ_D is an algebra morphism: $\chi_D \circ m_{\mathcal{H}_D^{\text{fg}}} = m_{\mathcal{H}^{\text{P}}} \circ (\chi_D \otimes \chi_D)$.

It is sufficient to prove this for the product of two generators $\Gamma_1, \Gamma_2 \in \mathcal{H}_D^{\text{fg}}$. Subdiagrams of the product $m(\Gamma_1 \otimes \Gamma_2) = \Gamma_1 \Gamma_2$, which is defined as the disjoint union of the two diagrams, can be represented as pairs (γ_1, γ_2) where $(\gamma_1, \gamma_2) \subset \Gamma_1 \Gamma_2$ if $\gamma_1 \subset \Gamma_1$ and $\gamma_2 \subset \Gamma_2$. This corresponds to the Cartesian product regarding the poset structure of the subdivergences. The loop number of such a pair is the sum of the loop numbers of the components. On those grounds:

$$\chi_D(\Gamma_1 \Gamma_2) = (\mathcal{P}_D^{\text{s.d.}}(\Gamma_1 \Gamma_2), \nu) = (\mathcal{P}_D^{\text{s.d.}}(\Gamma_1) \times \mathcal{P}_D^{\text{s.d.}}(\Gamma_2), \nu_1 + \nu_2) = \quad (45)$$

$$m_{\mathcal{H}^{\text{P}}}(\chi_D(\Gamma_1) \otimes \chi_D(\Gamma_2)). \quad (46)$$

To prove that χ_D is a coalgebra morphism, we need to verify that,

$$(\chi_D \otimes \chi_D) \circ \Delta_{\mathcal{H}_D^{\text{fg}}} = \Delta_{\mathcal{H}^{\text{P}}} \circ \chi_D. \quad (47)$$

Choosing some generator Γ of $\mathcal{H}_D^{\text{fg}}$ and using the definition of Δ_D :

$$(\chi_D \otimes \chi_D) \circ \Delta_{\mathcal{H}_D^{\text{fg}}} \Gamma = \sum_{\gamma \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)} \chi_D(\gamma) \otimes \chi_D(\Gamma/\gamma), \quad (48)$$

the statement follows from $\chi_D(\gamma) = ([\hat{0}, \gamma], \nu(\gamma))$ and $\chi_D(\Gamma/\gamma) = ([\emptyset, \Gamma/\gamma], \nu) \simeq ([\gamma, \Gamma], \nu - \nu(\gamma))$, which is a direct consequence of the definition of contractions in eq. (12). \square

Corollary 4. *In a join-meet-renormalizable QFT, $\text{im}(\chi_D) \subset \mathcal{H}^{\text{L}} \subset \mathcal{H}^{\text{P}}$, where \mathcal{H}^{L} is the subspace of \mathcal{H}^{P} which is generated by all elements (L, ν) , where L is a lattice. In other words: In a join-meet-renormalizable QFT, χ_D maps s.d. diagrams and products of them to decorated lattices.*

Example 8. For any primitive diagram $\Gamma \in \text{Prim}\mathcal{H}_D^{\text{fg}}$,

$$\chi_D(\Gamma) = (\mathcal{P}_D^{\text{s.d.}}(\Gamma), \nu) = \begin{array}{c} \textcircled{L} \\ | \\ \textcircled{0} \end{array}, \quad (49)$$

where the vertices in the Hasse-diagram are decorated by the value of ν and $L = h_1(\Gamma)$ is the loop number of the primitive diagram.

The coproduct of $\chi_D(\Gamma)$ in \mathcal{H}^{P} can be calculated using eq. (40):

$$\Delta_{\mathcal{H}^{\text{P}}} \begin{array}{c} \textcircled{L} \\ | \\ \textcircled{0} \end{array} = \begin{array}{c} \textcircled{L} \\ | \\ \textcircled{0} \end{array} \otimes \mathbb{I} + \mathbb{I} \otimes \begin{array}{c} \textcircled{L} \\ | \\ \textcircled{0} \end{array}. \quad (50)$$

As expected, these decorated posets are also primitive in \mathcal{H}^{P} .

Example 9. For the diagram $\times \circlearrowleft \times \in \mathcal{H}_4^{\text{fg}}$, χ_D gives the decorated poset,

$$\chi_D \left(\times \circlearrowleft \times \right) = \begin{array}{c} \textcircled{3} \\ \textcircled{2} \quad \textcircled{2} \\ \textcircled{1} \\ \textcircled{0} \end{array}, \quad (51)$$

of which the reduced coproduct in \mathcal{H}^P can be calculated,

$$\tilde{\Delta}_{\mathcal{H}^P} \begin{array}{c} \textcircled{3} \\ \textcircled{2} \quad \textcircled{2} \\ \textcircled{1} \\ \textcircled{0} \end{array} = 2 \begin{array}{c} \textcircled{2} \\ \textcircled{1} \\ \textcircled{0} \end{array} \otimes \begin{array}{c} \textcircled{1} \\ \textcircled{0} \end{array} + \begin{array}{c} \textcircled{1} \\ \textcircled{0} \end{array} \otimes \begin{array}{c} \textcircled{2} \\ \textcircled{1} \\ \textcircled{0} \end{array}. \quad (52)$$

This can be compared to the coproduct calculation in example 4,

$$\tilde{\Delta}_4 \times \circlearrowleft \times = 2 \times \circlearrowleft \times \otimes \times \circlearrowleft \times + \times \circlearrowleft \times \otimes \times \circlearrowleft \times \quad (53)$$

and identity eq. (47) is verified after computing the decorated poset of each subdiagram of $\times \circlearrowleft \times$ and comparing the previous two equations:

$$\chi_4 \left(\times \circlearrowleft \times \right) = \begin{array}{c} \textcircled{2} \\ \textcircled{1} \\ \textcircled{0} \end{array} \quad \chi_4 \left(\times \circlearrowleft \times \right) = \begin{array}{c} \textcircled{1} \\ \textcircled{0} \end{array} \quad \chi_4 \left(\times \circlearrowleft \times \right) = \begin{array}{c} \textcircled{2} \\ \textcircled{1} \\ \textcircled{0} \end{array}. \quad (54)$$

4 Properties of the lattices of subdivergences

Although, the above Hopf algebra morphism can be applied in every renormalizable QFT, we shall restrict ourselves to join-meet-renormalizable QFTs, where χ_D maps to \mathcal{H}^L , the Hopf algebra of decorated lattices, as a result of corollary 4.

The decorated lattice, which is associated to a Feynman diagram, encodes the ‘overlappingness’ of the diagrams’ subdivergences. Different join-meet-renormalizable QFTs have quite distinguished properties in this respect. Interestingly, the types of the decorated lattices appearing depend on the residues or equivalently on the superficial degree of divergence of the diagrams under consideration. For instance, it was proven by Berghoff in the context of Wonderful models that every diagram with only logarithmically divergent subdivergences (i.e. $\forall \gamma \in \mathcal{P}_D^{\text{s.d.}}(\Gamma) : \omega_D(\gamma) = 0$) is *distributive*:

Proposition 3. [2, Prop. 3.22] If Γ has only logarithmically s.d. subdiagrams in D dimensions, (i.e. $\forall \gamma \in \mathcal{P}_D^{\text{s.d.}}(\Gamma) \Rightarrow \omega_D(\gamma) = 0$), then the distributivity identities,

$$\gamma_1 \wedge (\gamma_2 \vee \gamma_3) = (\gamma_1 \wedge \gamma_2) \vee (\gamma_1 \wedge \gamma_3) \quad (55)$$

$$\gamma_1 \vee (\gamma_2 \wedge \gamma_3) = (\gamma_1 \vee \gamma_2) \wedge (\gamma_1 \vee \gamma_3), \quad (56)$$

hold for $\gamma_1, \gamma_2, \gamma_3 \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$.

This is a pleasant result for diagrams with only logarithmically divergent subdiagrams. Because distributive lattices are always graded, this implies that we have a bigrading on \mathcal{H}^L for these elements. One grading by the value of $\nu(\hat{1})$, corresponding to the loop number of the diagram, and one grading by the length of the maximal chains of the lattice, which coincides with the *coradical degree* of the diagram in $\mathcal{H}_D^{\text{fg}}$. The coradical filtration of $\mathcal{H}_D^{\text{fg}}$, defined in eq. (21), consequently becomes a grading for the subspaces generated by only logarithmically s.d. diagrams.

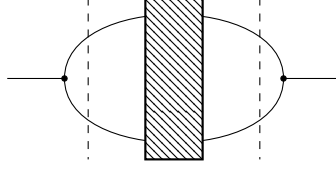


Figure 4: Structure of overlapping divergences in three-valent QFTs

4.1 Theories with only three-or-less-valent vertices

From the preceding result the question arises what structure is left, if we also allow subdiagrams which are not only logarithmically divergent. In renormalizable QFTs with only three-or-less-valent vertices, the lattices $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ will turn out to be semimodular. This is a weaker property than distributivity, but it still guarantees that the lattices are graded. To capture this property of $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$, some additional terms of lattice theory will be repeated following [14].

Join-irreducible element An element x of a lattice L , $x \in L$ is called join-irreducible if $x = y \vee z$ always implies $x = y$ or $x = z$.

Atoms and coatoms An element x of L is an atom of L if it covers $\hat{0}$. It is a coatom of L if $\hat{1}$ covers x .

Semimodular lattice A lattice L is semimodular if for two elements $x, y \in L$ that cover $x \wedge y$, x and y are covered by $x \vee y$.

With these notions we can formulate

Lemma 3. *In a QFT with only three-or-less-valent vertices, vertex-type s.d. diagrams Γ ($|H_{\text{ext}}(\Gamma)| = 3$) are always join-irreducible elements of $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$.*

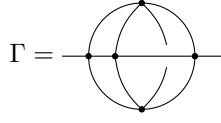
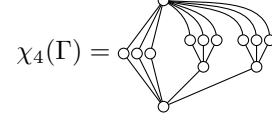
Proof. Suppose there were $\gamma_1, \gamma_2 \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$ with $\gamma_1 \neq \Gamma$, $\gamma_2 \neq \Gamma$ and $\gamma_1 \vee \gamma_2 = \Gamma$. The subdivergences γ_1 and γ_2 are therefore overlapping. From lemma 1 and proposition 1 we know $|H_{\text{ext}}(\gamma_1)| + |H_{\text{ext}}(\gamma_2)| = |H_{\text{ext}}(\gamma_1 \cup \gamma_2)| + |H_{\text{ext}}(\gamma_1 \cap \gamma_2)|$ and $|H_{\text{ext}}(\gamma_1 \cap \gamma_2)| \geq 4$. If $|H_{\text{ext}}(\gamma_1 \cup \gamma_2)| = 3$, these equations cannot be fulfilled for $|H_{\text{ext}}(\gamma_1)| \leq 3$ and $|H_{\text{ext}}(\gamma_2)| \leq 3$. \square

Proposition 4. In a renormalizable QFT with only three-or-less-valent vertices, the lattice $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ is semimodular for every Feynman diagram Γ .

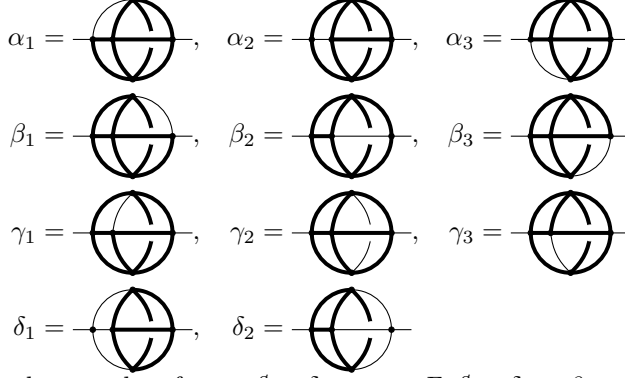
Proof. For two diagrams $\mu_1, \mu_2 \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$ we can always form the contractions by $\mu_1 \wedge \mu_2$: $\mu_1 / (\mu_1 \wedge \mu_2)$ and $\mu_2 / (\mu_1 \wedge \mu_2)$. Hence, the statement that μ_1, μ_2 cover $\mu_1 \wedge \mu_2$ is equivalent to stating that $\mu_1 / (\mu_1 \wedge \mu_2)$ and $\mu_2 / (\mu_1 \wedge \mu_2)$ are primitive.

To prove that $\mu_1 \vee \mu_2$ covers μ_1 and μ_2 if μ_1 and μ_2 cover $\mu_1 \wedge \mu_2$, it is therefore sufficient to verify that for γ_1, γ_2 primitive $(\gamma_1 \cup \gamma_2) / \gamma_1$ and $(\gamma_1 \cup \gamma_2) / \gamma_2$ are primitive as well. This is obvious if γ_1, γ_2 are not overlapping.

If γ_1 and γ_2 are overlapping, they must be of vertex-type and $\gamma_1 \cup \gamma_2$ of propagator-type. Because only three-valent vertices are allowed and $|H_{\text{ext}}(\gamma_i) \cap H_{\text{int}}(\gamma_1 \cup \gamma_2)| \geq 2$ (corollary 1), each γ_1 and γ_2 must provide one external edge for $\gamma_1 \cup \gamma_2$. The situation is depicted in fig. 4. For both γ_1 and γ_2 to be primitive, they must share the same four-leg kernel, depicted as a striped box. Contraction with either γ_1 or γ_2 results in a one loop propagator, which is primitive. \square

(a) Example of a diagram where $\mathcal{P}_4^{\text{s.d.}}(\Gamma)$ forms a non-graded lattice.

(b) The Hasse diagram of the corresponding non-graded lattice.



with the complete forests $\emptyset \subset \delta_1 \subset \alpha_i \subset \Gamma$, $\emptyset \subset \delta_2 \subset \beta_i \subset \Gamma$ and $\emptyset \subset \gamma_i \subset \Gamma$.

(c) The non-trivial superficially divergent subdiagrams and the complete forests which can be formed out of them.

Figure 5: Counter example of a lattice, which appears in join-meet-renormalizable QFTs with four-valent vertices and is not graded.

Semimodular lattices have a very rich structure, see for instance Stern's book [15]. Eventually, semimodularity implies that the lattices under consideration are graded:

Theorem 4. *In a renormalizable QFT with only three-or-less-valent vertices:*

- $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ is a graded lattice for every propagator, vertex-type diagram or disjoint unions of both.
- \mathcal{H}^L is bigraded by $\nu(\hat{1})$ and the length of the maximal chains of the lattices, which coincides with the coradical degree in \mathcal{H}^L .
- $\mathcal{H}_D^{\text{fg}}$ is bigraded by $h_1(\Gamma)$ and the coradical degree of Γ .
- Every complete forest of Γ has the same length.

Proof. Every semimodular lattice is graded [14, Proposition 3.3.2]. □

4.2 Theories with only four-or-less-valent vertices

We have shown that every lattice associated to a s.d. diagram in a QFT with only three-or-less-valent vertices is semimodular. For join-meet-renormalizable QFTs with also four-valent vertices the situation is more involved as the example in fig. 5 exposes. The depicted lattice in fig. 5b associated to the ϕ^4 -diagram in fig. 5a is not semimodular, because it is not graded. This implies that not all complete forests are of the same length in theories, where this topology can appear. This includes ϕ^4 and Yang-Mills theory in four dimensions. The s.d. subdiagrams of the counter

example are illustrated in fig. 5c. It can be seen that there are six complete forests of length four and three complete forests of length three.

The pleasant property of semimodularity can be recovered by working in the Hopf algebra of Feynman diagrams without tadpoles or equivalently by setting all tadpole diagrams to zero. This is quite surprising, because the independence of loops in tadpoles from external momenta and the combinatorial structure of BPHZ, encoded by the Hopf algebra of Feynman diagrams, seem independent on the first sight.

Formally, we can transfer the restriction to tadpole-free diagrams to \mathcal{H}^L by the following procedure: The Hopf algebra morphism $\psi : \mathcal{H}_D^{\text{fg}} \rightarrow \tilde{\mathcal{H}}_D^{\text{fg}}$ defined in eq. (26) gives rise to the ideal $\ker \psi \subset \mathcal{H}_D^{\text{fg}}$. Using the Hopf algebra morphism χ_D an ideal of \mathcal{H}^L , $\chi_D(\ker \psi) \subset \mathcal{H}^L$, is obtained. This can be summarized in a commutative diagram:

$$\begin{array}{ccc} \mathcal{H}^{\text{fg}} & \xrightarrow{\psi} & \tilde{\mathcal{H}}^{\text{fg}} \\ \downarrow \chi_D & & \downarrow \chi'_D \\ \mathcal{H}^L & \xrightarrow{\psi'} & \tilde{\mathcal{H}}^L \end{array}$$

where $\tilde{\mathcal{H}}^L$ is the quotient $\tilde{\mathcal{H}}^L := \mathcal{H}^L / \chi_D(\ker \psi)$ and ψ' is just the projection to $\tilde{\mathcal{H}}^L$.

The interesting part is the morphism $\chi'_D : \tilde{\mathcal{H}}_D^{\text{fg}} \rightarrow \tilde{\mathcal{H}}^L$, which maps from the Hopf algebra of Feynman diagrams without tadpoles to $\tilde{\mathcal{H}}^L$. Such a map can be constructed explicitly and for theories with only four-or-less-valent vertices, it can be ensured that χ'_D maps Feynman diagrams to decorated semimodular lattices.

Proposition 5. In a renormalizable QFT with only four-or-less-valent vertices, χ'_D maps elements from the Hopf algebra of Feynman diagrams without tadpoles to decorated lattices.

Proof. Explicitly, χ'_D is the map,

$$\chi'_D : \Gamma \mapsto (\tilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma), \nu), \quad (57)$$

where the decoration ν is the same as above.

We need to show that $\tilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma)$ ordered by inclusion is a lattice. This is not as simple as before, because $\gamma_1, \gamma_2 \in \tilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma)$ does not necessarily imply $\gamma_1 \cup \gamma_2 \in \tilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma)$. From definition 4 of $\tilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma)$, we can deduce that if $\gamma_1, \gamma_2 \in \tilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma)$, then $\gamma_1 \cup \gamma_2 \notin \tilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma)$ iff $\Gamma / \gamma_1 \cup \gamma_2$ is a tadpole.

To prove that there still exists a least upper bound for every pair γ_1, γ_2 we must ensure that every element $\mu \in \mathcal{P}_D^{\text{s.d.}}(\Gamma)$ and $\mu \notin \tilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma)$ is only covered by one element in $\tilde{\mathcal{P}}_D^{\text{s.d.}}(\Gamma)$. This is equivalent to stating that if $\gamma \subset \delta \subset \Gamma$ and δ/γ is a primitive tadpole (i.e. a self loop with one vertex), then there is no $\delta' \subset \Gamma$ such that δ'/γ is a primitive tadpole. There cannot be such a second subdiagram δ' . Suppose there were such δ and δ' . δ and δ' are obtained from γ by joining two of its external legs to an new edge. As only four-or-less-valent vertices are allowed, such a constellation can only be achieved if γ is a diagram with four external legs. δ and δ' are the diagrams obtained by closing either pair of legs of γ . This would imply that $\delta_1 \cup \delta_2$ is a vacuum diagram without external legs, which is excluded. \square

Example 10. The map χ'_D can be applied to the example in fig. 5 where the lattice obtained by χ_D is not semimodular. It can be seen from fig. 5c that the only diagrams, which do not

result in tadpole diagrams upon contraction are δ_1 and δ_2 . Accordingly,

$$\chi'_4 \left(\text{Diagram} \right) = \begin{matrix} \textcircled{5} \\ \textcircled{3} \textcircled{3} \\ \textcircled{0} \end{matrix}, \tag{58}$$

which is a graded lattice.

Proposition 6. In a QFT with only four-or-less-valent vertices χ'_D maps elements from the Hopf algebra of Feynman diagrams without tadpoles to decorated semimodular lattices.

Proof. As above we only need to prove that if γ_1 and γ_2 are overlapping and primitive, then $(\gamma_1 \vee \gamma_2)/\gamma_1$ and $(\gamma_1 \vee \gamma_2)/\gamma_2$ are primitive as well. The key is to notice that $H_{\text{ext}}(\gamma_1 \vee \gamma_2) \not\subset H_{\text{ext}}(\gamma_1)$ and $H_{\text{ext}}(\gamma_1 \vee \gamma_2) \not\subset H_{\text{ext}}(\gamma_2)$. Otherwise one of the contracted diagrams would be a tadpole. For this reason, we can characterize the overlapping divergences of a diagram by the proper subset of external half-edges of the full diagram it contains. If $\gamma_1 \vee \gamma_2/\gamma_1$ was not primitive, we could remove the half-edges $H_{\text{ext}}(\gamma_1) \cap H_{\text{ext}}(\gamma_1 \vee \gamma_2)$ and the adjacent vertices and edges. The result would be a s.d. subdiagram of γ_2 in contradiction with the requirement. \square

It is interesting how important taking the quotient by the tadpole diagrams is, to obtain the property of semimodularity for the lattices of Feynman diagrams.

Theorem 5. In a renormalizable QFT with only four-or-less-valent vertices:

- $\tilde{\mathcal{P}}_D^{s.d.}(\Gamma)$ is a graded lattice for every propagator, vertex-type diagram or disjoint unions of both.
- $\tilde{\mathcal{H}}^L$ is bigraded by $\nu(\hat{1})$ and the length of the maximal chains of the lattices, which coincides with the coradical degree in \mathcal{H}^L .
- $\tilde{\mathcal{H}}_D^{\text{fg}}$ is bigraded by $h_1(\Gamma)$ and the coradical degree of Γ .
- Every complete forest of Γ , which does not result in a tadpole upon contraction, has the same length.

Proof. Every semimodular lattice is graded [14, Proposition 3.3.2]. \square

The overlapping diagrams in $\tilde{\mathcal{P}}_D^{s.d.}(\Gamma)$ are characterized by the external legs of Γ they contain. As a consequence, there is a limited number of possibilities for primitive diagrams to be overlapping. A two-leg diagram can only be the join of at most two primitive overlapping diagrams and a three-leg diagram can only be the join of at most three primitive divergent overlapping diagrams. For four-leg diagrams in theories with only four-or-less valent vertices the restriction is even more serve: In these cases a four-leg diagram can only be the join of at most two primitive overlapping diagrams.

5 Applications to zero dimensional QFT and diagram counting

As an application of the lattice structure, the enumeration of some classes of primitive diagrams using techniques from zero-dimensional quantum field theories is presented. The characteristic property of zero-dimensional QFT is that every diagram in the perturbation expansion has the

amplitude 1 [8]. On the Hopf algebra of Feynman diagrams such a prescription can be formulated by the character or *Feynman rule*:

$$\phi : \mathcal{H}_D^{\text{fg}} \rightarrow \mathbb{Q}[\hbar], \quad (59)$$

$$\Gamma \mapsto \hbar^{h_1(\Gamma)}, \quad (60)$$

which maps every Feynman diagram to \hbar to the power of its number of loops in the ring of powerseries in \hbar . Clearly, ϕ is in $G_{\mathbb{Q}[\hbar]}^{\mathcal{H}_D^{\text{fg}}}$, the group of characters of $\mathcal{H}_D^{\text{fg}}$ to $\mathbb{Q}[\hbar]$. Note, that we are not setting $D = 0$ even though ϕ are the Feynman rules for zero-dimensional QFT. Every diagram would be ‘convergent’ and the Hopf algebra $\mathcal{H}_0^{\text{fg}}$ trivial. It might be more enlightening to think about ϕ as toy Feynman rules which assign 1 to every Feynman diagram without any respect to kinematics. This way we can still study the effects of renormalization on the amplitudes in an arbitrary dimension of space-time. Moreover, ϕ ‘counts’ the number of diagrams weighted by their symmetry factor.

We define the sum of all 1PI diagrams with a certain residue r weighted by their symmetry factor as,

$$X^r := \sum_{\substack{\Gamma \in \mathcal{T} \\ \text{res}(\Gamma)=r}} \frac{\Gamma}{|\text{Aut}(\Gamma)|}, \quad (61)$$

such that $\phi(X^r)$ is the generating function of these weighted diagrams with \hbar as a counting variable. This generating function is the perturbation expansion of the Green’s function for the residue r .

The *counter term map* [7] is defined as,

$$S_D^R := R \circ \phi \circ S_D, \quad (62)$$

in a *multiplicative renormalization scheme* R with the antipode S_D of $\mathcal{H}_D^{\text{fg}}$. S_D^R is called the counter term map, because it maps the sum of all 1PI diagrams with a certain residue r to the corresponding counter term, which when substituted into the lagrangian renormalizes the QFT appropriately. The renormalized Feynman rules are given by the convolution product $\phi_D^R := S_D^R * \phi$. In zero-dimensional QFTs, they vanish on all generators of $\mathcal{H}_D^{\text{fg}}$ except on $\mathbb{1}$. This can be used to obtain differential equations for the Z -factors and other interesting quantities as was done in [8, 1].

For the toy Feynman rules ϕ , there are no kinematics to choose a multiplicative renormalization scheme from. The renormalization will be modeled as usual in the scope of zero-dimensional-QFTs by setting $R = \text{id}$. Consequently, $S_D^R = \phi \circ S_D$.

The antipode in this formula is the point where the Hopf algebra structure enters the game. The lattice structure can be used to clarify the picture even more and to obtain quantitative results.

We define $\phi' \in G_{\mathbb{Q}[\hbar]}^{\mathcal{H}^L}$, a Feynman rule on the Hopf algebra of decorated lattices, analogous to $\phi \in G_{\mathbb{Q}[\hbar]}^{\mathcal{H}_D^{\text{fg}}}$:

$$\phi' : \mathcal{H}^L \rightarrow \mathbb{Q}[\hbar], \quad (63)$$

$$(P, \nu) \mapsto \hbar^{\nu(\hat{1})}, \quad (64)$$

which maps a decorated lattice to the value of the decoration of the largest element. Immediately, we can see that $\phi = \phi' \circ \chi_D$. For the counter term map,

$$S_D^R = \phi' \circ \chi_D \circ S_D, \quad (65)$$

is obtained. Using theorem 3, we can commute χ_D and S_D ,

$$S_D^R = \phi' \circ S \circ \chi_D. \quad (66)$$

For this reason, the evaluation of S_D^R can be performed entirely in \mathcal{H}^L ! S_D^R reduces to a combinatorial calculation on the lattice which is obtained by the Hopf algebra morphism χ_D . The morphism $\phi' \circ S$ maps decorated lattices into the ring of powerseries in \hbar . Because S respects the grading in $\nu(\hat{1})$, we can write

$$\phi' \circ S(L, \nu) = \hbar^{\nu(\hat{1})} \zeta \circ S(L, \nu), \quad (67)$$

where ζ is the characteristic function $(L, \nu) \mapsto 1$. It was shown in [9], that $\zeta \circ S$ is the *Moebius function*, $\mu(\hat{0}, \hat{1})$, on the lattice. It is defined recursively as,

Definition 10 (Moebius function).

$$\mu_P(x, y) = \begin{cases} 1, & \text{if } x = y \\ -\sum_{x \leq z < y} \mu_P(x, z) & \text{if } x < y. \end{cases} \quad (68)$$

for a poset P and $x, y \in P$.

We summarize these observations in

Theorem 6. *For zero-dimensional-QFT Feynman rules as ϕ above the counter term map takes the form*

$$S^{R'}(L, \nu) = \hbar^{\nu(\hat{1})} \mu_L(\hat{0}, \hat{1}) \quad (69)$$

on the Hopf algebra of lattices, where $S_D^R = S^{R'} \circ \chi_D$ and with $\hat{0}$ and $\hat{1}$ the lower and upper bound of L .

Corollary 5.

$$S_D^R(\Gamma) = \hbar^{h_1(\Gamma)} \mu_{\mathcal{P}_D^{s.d.}(\Gamma)}(\hat{0}, \hat{1}) \quad (70)$$

on the Hopf algebra of Feynman diagrams with $\hat{0} = \emptyset$ and $\hat{1} = \Gamma$, the lower and upper bound of $\mathcal{P}_D^{s.d.}(\Gamma)$.

On these grounds, the counter terms in zero-dimensional QFT can be calculated only by computing the Moebius function on the lattice $\mathcal{P}_D^{s.d.}(\Gamma)$! The Moebius function is a well studied object in combinatorics. There are especially sophisticated techniques to calculate the Moebius functions on lattices (see [14, 15]). The following two theorems, which are quoted from the respective sources, should serve as an illustration of this sophistication:

Theorem 7. *(Rota's crosscut theorem for atoms and coatoms (special case of [14, cor. 3.9.4]))*
Let L be a finite lattice and X its set of atoms and Y its set of coatoms, then

$$\mu_L(\hat{0}, \hat{1}) = \sum_k (-1)^k N_k = \sum_k (-1)^k M_k, \quad (71)$$

where N_k is the number of k -subsets of X whose join is $\hat{1}$ and M_k is the number of k -subsets of Y whose meet is $\hat{0}$.

With this theorem the Moebius functions of all the lattices appearing in this article can be calculated very efficiently.

In many cases, an even simpler theorem, which is a special case of the previous one, applies:

Theorem 8. (Hall's theorem [15, cor. 4.1.7.]) *If, in a lattice $\hat{1}$ is not a join of atoms or $\hat{0}$ is not a meet of coatoms, then $\mu(\hat{0}, \hat{1}) = 0$.*

In lemma 3, we proved that every vertex-type subdiagram in a QFT with only three-valent vertices is join-irreducible. Hence, it is also no join of atoms except if it is an atom itself.

Theorem 9. *In a renormalizable QFT with only three-or-less-valent vertices and Γ a vertex-type s.d. diagram (i.e. $H_{\text{ext}}(\Gamma) = 3$):*

$$S_D^R(\Gamma) = \begin{cases} -\hbar^{h_1(\Gamma)} & \text{if } \Gamma \text{ is primitive} \\ 0 & \text{if } \Gamma \text{ is not primitive.} \end{cases} \quad (72)$$

Proof. In both cases the element $\hat{1} = \Gamma$ in the lattice $\mathcal{P}_D^{\text{s.d.}}(\Gamma)$ is join-irreducible (lemma 3). If Γ is primitive $\phi \circ S(\Gamma) = -\phi(\Gamma) = -\hbar^{h_1(\Gamma)}$. If Γ is not primitive, it does not cover $\hat{0}$. This implies that $\hat{1}$ is not a join of atoms. Therefore, $\mu_{\mathcal{P}_D^{\text{s.d.}}(\Gamma)}(\hat{0}, \hat{1})$ vanishes and so does $S_D^R(\Gamma)$ in accordance to eq. (70). \square

Corollary 6. *In a renormalizable QFT with only three-or-less-valent vertices and $r \in \mathcal{R}_v$ a vertex-type residue:*

$$S_D^R(X^r) = -\phi \circ P_{\text{Prim}(\mathcal{H}_D^{\text{fg}})}(X^r), \quad (73)$$

where $P_{\text{Prim}(\mathcal{H}_D^{\text{fg}})}$ projects onto the primitive elements of $\mathcal{H}_D^{\text{fg}}$.

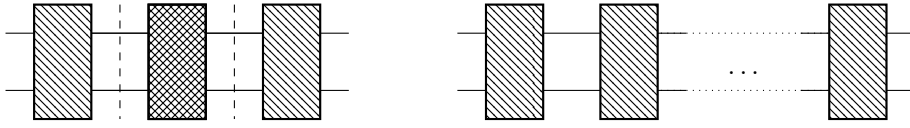
Summarizing, we established that in a theory with only three-or-less-valent vertices the counter term $S_D^R(X^r)$ counts the number of primitive diagrams if $r \in \mathcal{R}_v$. This fact has been used indirectly in [8] to obtain the generating functions for primitive vertex diagrams in ϕ^3 .

The conventional Z -factor for the respective vertex is $Z^r = 1 + r!S_D^R(X^r)$, where the factorial of the residue $r = \prod_{\varphi \in \Phi} \varphi^{n_\varphi}$ is $r! = \prod_{\varphi \in \Phi} n_\varphi!$. This factorial is necessary to return to the leg-fixed case.

Further exploitation of the lattice structure leads to a statement on propagator-type diagrams in such theories:

Theorem 10. *In a renormalizable QFT with three-or-less-valent vertices the propagator-type diagrams Γ , for which $S_D^R(\Gamma) \neq 0$, must have the lattice structure \circ or $\circ \circ$.*

Proof. A propagator-type diagram Γ either has a maximal forest which is the union of propagator diagrams, has at least two vertex-type subdiagrams or it is the primitive diagram of the topology $\circ \circ$. In the first case Γ is join-irreducible and $S_D^R(\Gamma) = 0$. In the third case the corresponding lattice is \circ . In the second case, Γ covers at least one vertex diagram γ which is join-irreducible. Every lattice L with $\mu_L(\hat{0}, \hat{1}) \neq 0$ is complemented [15, Cor. 4.1.11]. In a complemented lattice L , there is a $y \in L$ for every $x \in L$ such that $x \vee y = \hat{1}$ and $y \wedge x = \hat{0}$. For this reason, all the join-irreducible elements of L must be atoms if $\mu_L(\hat{0}, \hat{1}) \neq 0$. As was shown in the proof of proposition 4, a propagator cannot be the join of more than two primitive diagrams. Accordingly, $\circ \circ$ is the only possible lattice if Γ is not primitive. \square



(a) Structure of overlapping divergences of (b) Chain of overlapping divergences of type four-leg diagrams in theories with four-or-less- in fig. 6a, which will evaluate to a non-zero valent vertices. Moebius function.

Figure 6: Overlapping divergences for diagrams four legs in theories with only four-or-less-valent vertices.

The Z -factors for the propagators can also be obtained using the last theorem. To do this the Moebius function for each propagator diagram must be calculated using the form of the lattices and eq. (68). The Moebius functions for the vertex-type diagrams are known from theorem 9.

Example 11. In a renormalizable QFT with only three-or-less-valent vertices and $r \in \mathcal{R}_e$, a propagator-type residue:

$$S_D^R(X^r) = \hbar \sum_{\substack{\text{res } \Gamma_P = r \\ \{v_1, v_2\} = V(\Gamma_P)}} \frac{1}{\text{Aut } \Gamma_P} \left(-1 + \frac{\text{res}(v_1)!}{2} \phi \circ P_{\text{Prim}(\mathcal{H}_D^{\text{fg}})}(X^{\text{res}(v_1)}) + \frac{\text{res}(v_2)!}{2} \phi \circ P_{\text{Prim}(\mathcal{H}_D^{\text{fg}})}(X^{\text{res}(v_2)}) \right) \quad (74)$$

where the sum is over all primitive propagator diagrams Γ_P with a topology as \circlearrowleft and exactly two vertices v_1, v_2 . Of course, this sum is finite.

The factorials of the residues must be included to fix the external legs of the vertex-type subdiagrams. The factor of $\frac{1}{2}$ is necessary, because every non-primitive diagram, which contributes to the counter term, has exactly two maximal forests. This is an example of a simple Dyson-Schwinger equation in the style of [10].

The conventional Z -factor for the propagator is $Z^r = 1 + r! S_D^R(X^r)$.

Although the counter term map for renormalizable QFTs with only three-or-less valent vertices enumerates primitive diagrams, we cannot assume that the situation is similar in a more general setting with also four-valent vertices. A negative result in this direction was obtained in [1, p. 27] by Argyres et al. They observed that the vertex counter term in zero-dimensional ϕ^4 does not count primitive diagrams.

From the perspective of lattice theory this result can be explained. The only way in which overlapping divergences can appear in a diagram with four legs in a QFT with only four-or-less-valent vertices is depicted in fig. 6a. The dashed lines indicate the possible cuts to separate one overlapping divergence from the other. To obtain a Feynman diagram Γ with $\mu_{\mathcal{P}_D^{\text{s.d.}}(\Gamma)}(\hat{0}, \hat{1}) \neq 0$, the blob in the middle must be either of the same overlapping type as fig. 6a or superficially convergent. Otherwise, a join-irreducible element would be generated, which would imply $\mu_{\mathcal{P}_D^{\text{s.d.}}(\Gamma)}(\hat{0}, \hat{1}) = 0$. The possible non-primitive diagrams with four legs, which give a non-vanishing Moebius function are consequently of the form depicted in fig. 6b, where each blob must be replaced by a superficially convergent four-leg diagram such that the diagram remains 1PI. The lattice corresponding to this structure of overlapping divergences is a boolean lattice. Every superficially divergent subdiagram can be characterized by the particular set partition of ‘blocks’ it contains. This gives a bijection from $\{0, 1\}^n$ to all possible subdivergences. The Moebius function of boolean lattices evaluates to $(-1)^n$, where n is the number of atoms [14, Ex. 3.8.4.]. Accordingly, the structure of overlapping four-leg diagrams depends on the number of superficially convergent four-leg diagrams. The situation is especially simple in ϕ^4 -theory:

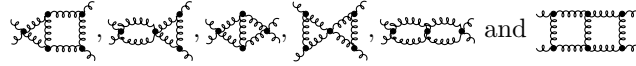
Example 12 (Overlapping vertex-type diagrams in ϕ^4 -theory). The only superficially convergent four-leg diagram in ϕ^4 -theory is the single four-leg vertex \times and so the only vertex-type s.d. diagrams Γ which give $\mu_{\mathcal{P}_D^{\text{s.d.}}(\Gamma)}(\hat{0}, \hat{1}) \neq 0$ are $\times \times \cdots \times$, the chains of one loop diagrams. Their generating function is $\frac{1}{8} \sum_{L \geq 0} \frac{\hbar^L}{2^L}$ (every diagram is weighted by its symmetry factor) and the counter term map in ϕ^4 evaluates to,

$$S_D^R(X^\times) = -\phi \circ P_{\text{Prim}(\mathcal{H}^{\text{fs}})}(X^\times) + \frac{1}{8} \sum_{L \geq 2} (-1)^L \left(\frac{\hbar}{2}\right)^L. \quad (75)$$

Note again, that in this setup the legs of the diagrams are not fixed as explained in section 2.2. To reobtain the numbers for the case with fixed legs, the generating function must be multiplied with the typical value $4! = 24$. The formula for the usual vertex Z -factor is $Z^\times = 1 + 4! S_D^R(X^\times)$.

Using this result, we can indeed use the counter terms of zero-dimensional ϕ^4 -theory to calculate the number of primitive diagrams. We merely must include the correction term on the right hand side of eq. (75).

Example 13 (Overlapping four-leg-vertex diagrams in pure Yang-Mills-theory). In pure Yang-Mills-theory there can be either the single four-valent vertex \bowtie or two three-valent vertices joined by a propagator --- as superficially convergent four-leg diagrams. Only chains of diagrams as in fig. 6b or primitive diagrams give a non-zero $S_D^R(\Gamma)$. At two loop for instance, the non-primitive diagrams



contribute non-trivially to $S_D^R(X^{\bowtie})$. These are the only four-leg diagrams which can be formed as the union of two primitive diagrams in this theory. The generating function for $L \geq 2$ of these diagrams is $\frac{3}{8} \sum_{L \geq 2} \left(\frac{3\hbar}{2}\right)^L$. Hence, the counter term map in pure Yang-Mills-theory for the four-gluon amplitude in zero-dimensional QFT evaluates to,

$$S_D^R(X^{\bowtie}) = -\phi \circ P_{\text{Prim}(\mathcal{H}^{\text{fs}})}(X^{\bowtie}) + \frac{3}{8} \sum_{L \geq 2} (-1)^L \left(\frac{3\hbar}{2}\right)^L. \quad (76)$$

To reobtain the numbers for the case with fixed legs this generating function needs to be multiplied with the value $4! = 24$ as in the example for ϕ^4 -theory. The formula for the Z -factor is $Z^{\bowtie} = 1 + 4! S_D^R(X^{\bowtie})$.

6 Outlook

It was shown that the Hopf algebra morphism χ_D can be used to disentangle the combinatorial and the analytic part of Feynman diagram calculations. The Hopf algebra of Feynman diagrams is graded by the coradical degree of its elements as demonstrated in theorem 4 for the case of three-or-less-valent theories and theorem 5 for four-or-less-valent theories with the tadpole diagrams set to zero.

This grading could also be used to reorganize diagrams which ‘renormalize in the same way’. The preimage χ_D^{-1} of some a decorated lattice in $\tilde{\mathcal{H}}^L$ is a space of such diagrams. With methods from [5] this could be used to express the log-expansion of Green functions systematically.

Primitive diagrams of coradical degree one contribute to the first power in the log-expansion, diagrams of coradical degree two to the second and diagrams with a coradical degree equal to the loop number contribute to the leading-log [11].

Furthermore, this framework can be used to make more statements and explicit calculations on the weighted numbers of primitive diagrams in different QFTs and their asymptotic behaviour. These aspects will be analyzed from a combinatorial perspective in a subsequent article [4].

Acknowledgements

I thank Dirk Kreimer for his great supervision, full support and encouragement, David Broadhurst for sparking my interest in the subject of zero-dimensional QFT, for many helpful conversations and motivation to write these results down, Erik Panzer for discussions and for finding the diagram in fig. 3, with a set of subdivergences which does not form a lattice, as well as Dominique Manchon and Marko Berghoff for fruitful discussions on the subject.

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