# A UNIFYING THEORY OF A POSTERIORI FINITE ELEMENT ERROR CONTROL 

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#### Abstract

Residual-based a posteriori error estimates are derived within a unified setting for lowest-order conforming, nonconforming, and mixed finite element schemes. The various residuals are identified for all techniques and problems as the operator norm $\|\ell\|$ of a linear functional of the form $$
\ell(v):=\int_{\Omega} p_{h}: D v d x+\int_{\Omega} g_{\Omega} \cdot v d x
$$ in the variable $v$ of a Sobolev space $V$. The main assumption is that the first-order finite element space $\mathcal{S}_{0}^{1}(\Omega) \subset \operatorname{ker} \ell \subset V$ is included in the kernel ker $\ell$ of $\ell$. As a consequence, any residual estimator that is a computable bound of $\|\ell\|$ can be used within the proposed frame without further analysis for nonconforming or mixed FE schemes. Applications are given for the Laplace, Stokes, and Navier-Lamè equations.


## 1. Unifying Theory of A Posteriori Error Control

This section sets up an abstract framework for a posteriori estimation which is filled with details for low-order finite element methods for the Laplace, Stokes, and Navier-Lamè equations in Section 2, 3, and 4, respectively. This unifying approach generalizes known techniques based on a Helmholtz decomposition [A, C, CD, CF1, CBK, BC, DDVP] as well as comparison schemes [HW]. The final result of the presented theory is that one unified type of residuals has to be analyzed once and the resulting estimator can be simultaneously used for a posteriori error control of conforming, nonconforming, mixed, or other finite element schemes.
1.1. Residual-Based A Posteriori Error Control. Let $A: X \rightarrow$ $X^{*}$ be a linear and bounded operator between the (real) Banach spaces $X$ and its dual $X^{*}$. Suppose that $A$ is surjective and injective such

[^0]that $A^{-1}$ is bounded as well. In particular, given $y \in X^{*}$, there exists a unique $x \in X$ such that
\[

$$
\begin{equation*}
A x=y \tag{1.1}
\end{equation*}
$$

\]

Suppose we are given some finite element approximation $x_{h} \in X$ (possibly with some additional properties generated by the computational scheme that provided $x_{h}$ ). Then we address the issue of approximating the error $\left\|x-x_{h}\right\|$. Notice that the linear functional

$$
\begin{equation*}
\mathcal{R e s}:=y-A x_{h}=A\left(x-x_{h}\right) \in X^{*} \tag{1.2}
\end{equation*}
$$

is known or, at least computable. Throughout this paper, an inequality $a \lesssim b$ replaces $a \leq c b$ with a multiplicative mesh-size independent constant $c$ that depends only on the domain $\Omega$ and the shape (e.g. through the aspect ratio) of finite elements. Finally, $a \approx b$ abbreviates $a \lesssim b \lesssim a$. Since the operator norms $\|A\| \approx 1 \approx\left\|A^{-1}\right\|$ of $A$ and $A^{-1}$ are uniformly bounded, there holds

$$
\begin{equation*}
\left\|x-x_{h}\right\|_{X} \approx\left\|A\left(x-x_{h}\right)\right\|_{X^{*}} \approx\|\mathcal{R} e s\|_{X^{*}} \tag{1.3}
\end{equation*}
$$

Hence, any residual-based a posteriori error control means the approximation of lower and upper bounds of the dual norm $\|\mathcal{R} e s\|_{X^{*}}$ of $\mathcal{R e s}$. Throughout the paper, $X=H \times L$ will (essentially) be fixed and the discrete subspaces vary.
1.2. Goal-Oriented Error Control. The analysis of this paper focuses on the estimation of the norm $\left\|x-x_{h}\right\|_{X}$. In some applications, there is a given (hence known) linear and bounded functional $\rho: X \rightarrow \mathbb{R}$ that monitors the error $\left|\rho\left(x-x_{h}\right)\right|$ (e.g. the error of an averaged strain or traction over a small but fixed region). To assess the latter quantity, let $A^{*}: X \rightarrow X^{*}$ be the dual operator of $A$ (for reflexive spaces $X=X^{* *}$ ) and let $z \in X$ be the solution to

$$
A^{*} z=\rho .
$$

Then it remains to estimate

$$
\rho\left(x-x_{h}\right)=\left(A^{*} z\right)\left(x-x_{h}\right)=\left(A\left(x-x_{h}\right)\right) z=\mathcal{R} \operatorname{es}(z) .
$$

An immediate consequence of this reads

$$
\begin{equation*}
\left|\rho\left(x-x_{h}\right)\right| \leq\|\mathcal{R} e s\|_{X^{*}}\|z\|_{X} \approx\|\rho\|_{X^{*}}\|\mathcal{R} e s\|_{X^{*}} \tag{1.4}
\end{equation*}
$$

This global estimate (1.4) is (a) presumably too coarse and (b) does not convey local information of $\rho$ via $z$. But it indicates that the evaluation of $\mathcal{R} \operatorname{es}(z)$ may follow localized arguments from the assessment of $\|\mathcal{R} e s\|_{X^{*}}$ addressed in this paper (cf. [BR1, AO, BaS]).
1.3. Mixed Approach to Flux or Stress Error Control. The primal variable $u \in H$ (e.g. the displacement field) is accompanied by a dual variable $p \in L$ (e.g. the flux or stress). The pair $(p, u)=: x \in$ $L \times H=: X$ plays the role of the variable $x$ in Subsection 1.1 above. Below, $L$ will be a Lebesgue space (e.g. $\left.L=L^{2}(\Omega)^{n}\right)$ and $H$ will be a Sobolev space (e.g. $H=H_{0}^{1}(\Omega)$ ), defined on a bounded domain $\Omega$ in $\mathbb{R}^{n}$. At the moment, it suffices to consider $L$ and $H$ as reflexive Banach spaces. The linear operator $A: X \rightarrow X^{*}$ is defined via a mixed framework, namely,

$$
\begin{equation*}
(A(p, u))(q, v):=a(p, q)+b(p, v)+b(q, u) \tag{1.5}
\end{equation*}
$$

for bounded bilinear forms $a: L \times L \rightarrow \mathbb{R}$ and $b: L \times H \rightarrow \mathbb{R}$. Under well-analyzed conditions on $a$ and $b[\mathrm{~B}, \mathrm{BF}]$ the operator $A$ is bijective. Hence, given right-hand sides $f \in L^{*}$ and $g \in H^{*}$ with $y \in X^{*}$ defined by $y(q, v)=f(q)+g(v)$, there exists a unique $x=(p, u) \in L \times H$ that solves (1.1). Let $x_{h}=\left(p_{h}, \tilde{u}_{h}\right) \in L \times H$ be an approximation to $x$ and define $\mathcal{R} e s$ as $\mathcal{R} e s_{L}+\mathcal{R} e s_{H}$ by (1.2), namely

$$
\begin{align*}
\operatorname{Res}_{L}(q) & :=f(q)-a\left(p_{h}, q\right)-b\left(q, \tilde{u}_{h}\right) \quad \text { for } q \in L, \\
\operatorname{Res}_{H}(v) & :=g(v)-b\left(p_{h}, v\right) \text { for } v \in H . \tag{1.6}
\end{align*}
$$

The notation $\tilde{u}_{h} \in H$ here and below asserts that $\tilde{u}_{h}$ is a continuous and not necessarily a discrete function; the subindex in $\tilde{u}_{h}$ refers to the fact that $\tilde{u}_{h}$ is closely related to $u_{h}$ and is on our disposal. With (1.5)-(1.6), Equivalence (1.3) becomes

$$
\begin{equation*}
\left\|p-p_{h}\right\|_{L}+\left\|u-\tilde{u}_{h}\right\|_{H} \approx\left\|\mathcal{R} e s_{L}\right\|_{L^{*}}+\left\|\operatorname{Res}_{H}\right\|_{H^{*}} . \tag{1.7}
\end{equation*}
$$

This is the starting point of the unifying theory. The fact that $\tilde{u}_{h}$ has to belong to $H$ is a crucial point in the sequel.

Remark 1.1. For non-conforming or mixed finite element schemes we obtain an approximation $u_{h}$ to $u$ which, below, is not in $H$. Consequently, $\tilde{u}_{h} \in H$ is, in general, different from $u_{h}$. To achieve an error estimation of the dual variable $p-p_{h}$, we will choose $\tilde{u}_{h}$ properly. The choice $\tilde{u}_{h}=u$ might be possible and minimizes $\left\|u-\tilde{u}_{h}\right\|$ but, in general, leads to difficulties in the evaluation of $b\left(q, \tilde{u}_{h}\right)$ in $\mathcal{R} e s_{L}$ in $(1.6)_{a}$.

Remark 1.2. It should be notified clearly that (1.5) is a primal mixed formulation, also called hybrid in $[\mathrm{BF}]$, where $L$ is not a subspace of $H(\operatorname{div} ; \Omega)$. This is because the derivatives act on $u \in H=H_{0}^{1}(\Omega)^{m}$. For instance, the bilinear form $b$ looks like

$$
b(q, v)=-\int_{\Omega} q \cdot \nabla v d x \quad \text { for the Laplace equation }
$$

(without an integration by parts as for the dual mixed formulation) and similar expressions hold for Stokes and Navier-Lamè equations.
1.4. Residuals. It is an aim of this paper to emphasize that there is essentially only one type of residual that arise in a posteriori error control: Given some $g \in L^{2}(\Omega)^{m}$ and some $g_{\mathcal{E}} \in L^{2}(\cup \mathcal{E})^{m}$ one encounters the linear functional Res : $H_{0}^{1}(\Omega)^{m} \rightarrow \mathbb{R}$,

$$
\begin{equation*}
\operatorname{Res}(v):=\int_{\Omega} g \cdot v d x+\int_{\cup \mathcal{E}} g_{\mathcal{E}} \cdot v d s \tag{1.8}
\end{equation*}
$$

with the property that

$$
\begin{equation*}
\mathcal{S}_{0}^{1}(\mathcal{T})^{m} \subseteq \operatorname{ker} \operatorname{Res} \tag{1.9}
\end{equation*}
$$

for globally continuous and piecewise polynomials $\mathcal{S}^{1}(\mathcal{T})^{m}$ of (partial or total) degree $\leq 1$ and $\mathcal{S}_{0}^{1}(\mathcal{T}):=\mathcal{S}^{1}(\mathcal{T}) \cap H_{0}^{1}(\Omega)$. (Details on the notation follow in Section 2). The norm of Res enters as an explicit bound in the form of

$$
\|\operatorname{Res}\|:=\sup _{v \in H_{0}^{1}(\Omega)^{m} \backslash\{0\}} \operatorname{Res}(v) /\|\nabla v\|_{L^{2}(\Omega)}
$$

or, for $m=3=n$ and nonconforming terms, in the form

$$
\|\operatorname{Res}\|:=\sup _{v \in H^{1}(\Omega)^{3}, \operatorname{Curl} v \not \equiv 0} \operatorname{Res}(v) /\|\operatorname{Curl} v\|_{L^{2}(\Omega)}
$$

It will be seen in Section 2, 3, and 4 that all arising residuals can be written and hence estimated in this unified form (1.8)-(1.9).

## 2. Application to Laplace Equation

This section is devoted to the Poisson problem as the simplest elliptic PDE and its residual-based a posteriori finite element error control. Subsection 2.1 introduces the model problem and some required notation while Subsection $2.2,2.3$, and 2.4 concern technical details in increasing difficulty for the conforming, non-conforming, and mixed loworder finite element methods. An application to discontinuous Galerkin schemes is in preparation.
2.1. Model Problem. Throughout this paper, $\Omega$ denotes a bounded Lipschitz domain in $\mathbb{R}^{n}$ with piecewise flat boundary $\partial \Omega$ such that $\bar{\Omega}$ is the union of a regular triangulation $\mathcal{T}, \bar{\Omega}=\cup \mathcal{T}$ (no hanging nodes). The Lebesgue and Sobolev spaces $L^{2}(\Omega)$ and $H^{1}(\Omega)$ are defined as usual and we define

$$
\begin{equation*}
L:=L^{2}(\Omega)^{n} \text { and } H:=H_{0}^{1}(\Omega):=\left\{w \in H^{1}(\Omega): w=0 \text { on } \partial \Omega\right\} \tag{2.1}
\end{equation*}
$$

Then, the gradient operator, $\nabla: H \rightarrow L$, maps $H$ into $L$. Given $g \in L^{2}(\Omega)$ let $u \in H$ denote the solution to the Poisson Problem

$$
\begin{equation*}
\Delta u+g=0 \text { in } \Omega \quad \text { and } \quad u=0 \text { on } \partial \Omega . \tag{2.2}
\end{equation*}
$$

Then, the flux $p:=\nabla u \in L$ and $u \in H$ solve the problem

$$
\begin{align*}
& (A(p, u))(q, v):=a(p, q)+b(p, v)+b(q, u) \\
& \quad \stackrel{!}{=}-\int_{\Omega} g v d x \quad \text { for all }(q, v) \in X=L \times H \tag{2.3}
\end{align*}
$$

of the form considered in Subsection 1.3 with

$$
\begin{equation*}
a(p, q):=\int_{\Omega} p \cdot q d x \quad \text { and } \quad b(p, v):=-\int_{\Omega} p \cdot \nabla v d x \tag{2.4}
\end{equation*}
$$

Theorem 2.1. The operator $A: X \rightarrow X^{*}$ defined in (2.3) a is bounded, linear, and bijective. For any $p_{h} \in L$ and $\tilde{u}_{h} \in H$ there holds

$$
\begin{equation*}
\left\|p-p_{h}\right\|_{L}+\left\|u-\tilde{u}_{h}\right\|_{H} \approx\left\|\mathcal{R} e s_{L}\right\|_{L^{*}}+\left\|\mathcal{R} e s_{H}\right\|_{H^{*}} \tag{2.5}
\end{equation*}
$$

for $\mathcal{R e s}_{L} \in L^{*}$ and $\mathcal{R e s}_{H} \in H^{*}$ defined for all $q \in L$ and $v \in H$ by

$$
\begin{align*}
\mathcal{R} e s_{L}(q) & :=\int_{\Omega} q \cdot\left(\nabla \tilde{u}_{h}-p_{h}\right) d x \quad \text { and }  \tag{2.6}\\
\mathcal{R} e s_{H}(v) & :=-\int_{\Omega} g v d x+\int_{\Omega} p_{h} \cdot \nabla v d x .
\end{align*}
$$

Proof. The assertions on $A$ are well known; a direct proof of an inf-sup condition follows for any $(p, u) \in L \times H$ from $(q, v):=(p-\nabla u ;-2 u) \in$ $L \times H$ and

$$
\begin{aligned}
1 / 6\|(p, u)\|_{X}\|(q, v)\|_{X} & \leq 1 / 6\left(\|p\|_{L}+\|u\|_{H}\right)\left(\|p\|_{L}+3\|u\|_{H}\right) \\
& \leq\|p\|_{L}^{2}+\|u\|_{H}^{2}=(A(p, u))(q, v) .
\end{aligned}
$$

The (generalised) Lax-Milgram lemma then yields bijectivity of $A$. The remaining assertions follow with the arguments of Subsection 1.3 which lead to (1.7) which, here, reads (2.5)-(2.6).
2.2. Conforming Finite Element Methods. The aforementioned triangulation $\mathcal{T}$ into triangles or parallelograms for 2D and into tetrahedra or parallelepipeds for 3D is the basis of the conforming loworder finite element space $H_{h}=\mathcal{S}_{0}^{1}(\mathcal{T})$. Let $\mathcal{P}_{k}(T)=P_{k}(T)$ and $\mathcal{P}_{k}(T)=Q_{k}(T)$ for a triangle (or tetrahedron) and parallelogram (or parallelepiped), respectively, and the space $P_{k}(T)$ and $Q_{k}(T)$ algebraic polynomials of total and partial degree $\leq k$, respectively, and define

$$
\begin{align*}
& \mathcal{L}^{k}(\mathcal{T}):=\left\{v \in L^{2}(\Omega): \forall T \in \mathcal{T},\left.v\right|_{T} \in \mathcal{P}_{k}(T)\right\} \quad \text { for } k=0,1 \\
& \mathcal{S}^{1}(\mathcal{T}):=\mathcal{L}^{1}(\mathcal{T}) \cap C(\bar{\Omega}) \quad \text { and } \quad \mathcal{S}_{0}^{1}(\mathcal{T}):=\mathcal{S}^{1}(\mathcal{T}) \cap H_{0}^{1}(\Omega) \tag{2.7}
\end{align*}
$$

Let $\mathcal{N}$ denote the set of nodes (i.e. the vertices of elements in $\mathcal{T}$ ) and let $\mathcal{E}$ denote the edges in 2D (or faces in 3D) in $\mathcal{T}$. Let $h_{\mathcal{T}}$ and $h_{\mathcal{E}}$ be $\mathcal{T}$ - and $\mathcal{E}$-piecewise constant on $\Omega$ and $\cup \mathcal{E}=\cup_{E \in \mathcal{E}} E$ defined by $\left.h_{\mathcal{T}}\right|_{T}:=h_{T}:=\operatorname{diam}(T)$ and $\left.h_{\mathcal{E}}\right|_{E}:=h_{E}:=\operatorname{diam}(E)$ for $T \in \mathcal{T}$ and $E \in \mathcal{E}$, respectively.

Given $u_{h} \in H_{h}$ with

$$
\begin{equation*}
\int_{\Omega} \nabla u_{h} \cdot \nabla v_{h} d x=\int_{\Omega} g v_{h} d x \quad \text { for all } v_{h} \in H_{h} \tag{2.8}
\end{equation*}
$$

and $p_{h}:=\nabla u_{h}$ as an approximation to $p:=\nabla u$, we aim to estimate $\left\|p-p_{h}\right\|_{L}$. With Theorem 2.1 and $\tilde{u}_{h}=u_{h} \in H_{h} \subset H$ there holds

$$
\begin{equation*}
\mathcal{R} e s_{L} \equiv 0 \quad \text { and } \quad \mathcal{R} e s_{H}(v)=-\int_{\Omega} g v d x+\int_{\Omega} p_{h} \cdot \nabla v d x \tag{2.9}
\end{equation*}
$$

Notice that an elementwise integration by parts shows that $\mathcal{R} e s_{H}$ is of the form (1.8)-(1.9). The evaluation of the residual $\mathcal{R} e s_{H}$, namely the estimation of lower and upper bounds of

$$
\begin{equation*}
\left\|\mathcal{R} e s_{H}\right\|_{H^{*}}:=\sup _{v \in H \backslash\{0\}}\left(\int_{\Omega} g v d x-\int_{\Omega} p_{h} \cdot \nabla v d x\right) /\|v\|_{H}, \tag{2.10}
\end{equation*}
$$

is subject of a vast literature. Although possibly sometimes not stated explicitly in this form, it is in fact the content of the books [V2, EJ, $\mathrm{AO}, \mathrm{BaS}]$. The point in this paper is that any of the (energy error) estimators thereof can be used. The standard explicit estimator reads

$$
\begin{equation*}
\eta_{R}^{(1)}:=\left\|h_{\mathcal{T}}\left(g+\operatorname{div}_{\mathcal{T}} p_{h}\right)\right\|_{L^{2}(\Omega)}+\left\|h_{\mathcal{E}}^{1 / 2}\left[p_{h} \cdot \nu_{\mathcal{E}}\right]\right\|_{L^{2}\left(\cup \mathcal{E}_{\Omega}\right)} \tag{2.11}
\end{equation*}
$$

and can be refined [CV] for $g \in H^{1}(\Omega)$ to

$$
\begin{equation*}
\eta_{R}^{(2)}:=\left\|h_{\mathcal{T}}^{2} \nabla g\right\|_{L^{2}(\Omega)}+\left\|h_{\mathcal{E}}^{1 / 2}\left[p_{h} \cdot \nu_{\mathcal{E}}\right]\right\|_{L^{2}\left(\cup \mathcal{E}_{\Omega}\right)} \tag{2.12}
\end{equation*}
$$

Another simple and easy-to-evaluate estimate is based on gradientrecovery: For any node $z \in \mathcal{N}$ with patch $\omega_{z}:=\operatorname{int}(\cup\{T \in \mathcal{T}: z \in T\})$ let $A_{z} p_{h}:=\int_{\omega_{z}} p_{h} d x /\left|\omega_{z}\right| \in \mathbb{R}^{n}$ be the average of $p_{h}$ on $\omega_{z}$. With the nodal basis function $\varphi_{z}$ (defined by $\varphi_{z} \in \mathcal{S}^{1}(\mathcal{T})$ and $\varphi_{z}(z)=1$ and $\varphi_{z}(x)=0$ for all $\left.x \in \mathcal{N} \backslash\{z\}\right)$ let

$$
\begin{equation*}
\eta_{A}:=\left\|h_{\mathcal{T}}^{2} \nabla g\right\|_{L^{2}(\Omega)}+\left\|p_{h}-\sum_{z \in \mathcal{N}}\left(A_{z} p_{h}\right) \varphi_{z}\right\|_{L^{2}(\Omega)} \tag{2.13}
\end{equation*}
$$

There holds (assuming $g \in H^{1}(\Omega)$ and $\left\|h_{\mathcal{T}}^{2} \nabla g\right\|_{L^{2}(\Omega)} \lesssim\left\|\mathcal{R} e s_{H}\right\|_{H^{*}}$ )

$$
\begin{equation*}
\left\|\mathcal{R e} s_{H}\right\|_{H^{*}} \approx \eta_{R}^{(1)} \approx \eta_{R}^{(2)} \approx \eta_{A} . \tag{2.14}
\end{equation*}
$$

Remark 2.1. Notice that $\left\|h_{\mathcal{T}}^{2} \nabla g\right\|_{L^{2}(\Omega)}$ is of higher order for the loworder finite element scheme analyzed in this paper.

Remark 2.2. The proof of $\left\|p-p_{h}\right\|_{L} \lesssim \eta_{R}^{(1)}$ goes back to [BaR, BaM, EJ], the proof of $\eta_{R}^{(1)} \lesssim\left\|p-p_{h}\right\|_{L}+\left\|h_{\mathcal{T}}^{2} \nabla g\right\|_{L^{2}(\Omega)}$ to [BaM, V2].

Remark 2.3. There are more expensive implicit error estimates, cf. [AO, V2, BaS, CBK].
2.3. Nonconforming Finite Element Methods. Based on the regular triangulation $\mathcal{T}$ into simplices (no parallelograms), the non-conforming finite element schemes due to Crouzeix-Raviart reads

$$
\begin{array}{r}
\mathcal{S}_{0}^{1, N C}(\mathcal{T}):=\left\{v \in \mathcal{L}^{1}(\mathcal{T}): v \text { continuous at } \mathcal{M} \cap \Omega\right.  \tag{2.15}\\
\text { and } v=0 \text { at } \mathcal{M} \cap \partial \Omega\}
\end{array}
$$

where $\mathcal{M}$ is the set of midpoints of edges (of faces) $E \in \mathcal{E}$. Notice that

$$
\mathcal{S}_{0}^{1}(\mathcal{T}) \subset \mathcal{S}_{0}^{1, N C}(\mathcal{T}) \subset H^{1}(\mathcal{T}):=\left\{v \in L^{2}(\Omega): \forall T \in \mathcal{T},\left.v\right|_{T} \in H^{1}(T)\right\}
$$

and this is, in general, $\not \subset H^{1}(\Omega)$. Let $\nabla_{\mathcal{T}}$ (resp. $\nabla_{\mathcal{T}}^{k}$ ) denote the $\mathcal{T}$ piecewise action of the gradient operator (resp. the matrix of all partial derivatives of order $k)$. Then, for any $v_{h} \in \mathcal{S}_{0}^{1, N C}(\mathcal{T}), \nabla_{\mathcal{T}} v_{h} \in L$. The finite element solution $u_{h} \in \mathcal{S}_{0}^{1, N C}(\mathcal{T})$ is the unique solution to

$$
\begin{equation*}
\int_{\Omega} \nabla_{\mathcal{T}} u_{h} \cdot \nabla_{\mathcal{T}} v_{h} d x=\int_{\Omega} g v_{h} d x \quad \text { for all } v_{h} \in \mathcal{S}_{0}^{1, N C}(\mathcal{T}) . \tag{2.16}
\end{equation*}
$$

The aim is to estimate the flux error $p-p_{h}$ for the discrete flux $p_{h}:=$ $\nabla_{\mathcal{T}} u_{h} \in L$. One difficulty is that, in general, $u_{h} \notin H$ and so $\tilde{u}_{h}$ cannot be chosen as in Subsection 2.2. However, for any $\tilde{u}_{h} \in H$, Theorem 2.1 yields

$$
\begin{equation*}
\left\|p-p_{h}\right\|_{L}+\left\|u-\tilde{u}_{h}\right\|_{H} \approx\left\|\mathcal{R} e s_{L}\right\|_{L^{*}}+\left\|\mathcal{R} e s_{H}\right\|_{H^{*}} \tag{2.17}
\end{equation*}
$$

with $\mathcal{R e s}_{H}$ treated as in (2.9)-(2.14) (notice that (2.16) guarantees (1.9) because of $\mathcal{S}_{0}^{1}(\mathcal{T}) \subset \mathcal{S}_{0}^{1, N C}(\mathcal{T})$ ) and

$$
\begin{equation*}
\mathcal{R e s}_{L}(q):=\int_{\Omega}\left(\nabla \tilde{u}_{h}-p_{h}\right) \cdot q d x \quad \text { for all } q \in L \tag{2.18}
\end{equation*}
$$

The focus of the remaining part of this subsection is therefore on

$$
\begin{equation*}
\left\|\mathcal{R} e s_{L}\right\|_{L^{*}}=\sup _{q \in L \backslash\{0\}} \operatorname{Res}_{L}(q) /\|q\|_{L}=\left\|p_{h}-\nabla \tilde{u}_{h}\right\|_{L^{2}(\Omega)} \tag{2.19}
\end{equation*}
$$

and so on a proper choice of $\tilde{u}_{h}$.
Remark 2.4. A direct approach towards an upper bound of (2.19) is to compute some $\tilde{u}_{h}:=I u_{h} \in \mathcal{S}_{0}^{1}(\mathcal{T})$ from the nonconforming finite element solution $u_{h}$; cf. [HW] for an example of $I$.

Remark 2.5. The minimizing $\tilde{u}_{h}$ in (2.19) within the class of conforming finite element approximations is equal to the conforming finite element approximation $u_{h}^{c}$ of Subsection 2.2. [Proof: Since $\int_{\Omega} p_{h} \cdot \nabla v_{h}^{c} d x=$ $\int_{\Omega} g v_{h}^{c} d x=\int_{\Omega} \nabla u_{h}^{c} \cdot \nabla v_{h}^{c} d x$ for all $v_{h}^{c} \mathcal{S}_{0}^{1}(\mathcal{T}) \subset \mathcal{S}_{0}^{1, N C}(\mathcal{T})$ there holds $\int_{\Omega}\left(p_{h}-\nabla u_{h}^{c}\right) \cdot \nabla v_{h}^{c} d x=0$.] The split $\nabla_{\mathcal{T}} u_{h}=\nabla u_{h}^{c}+\operatorname{Curl}_{\mathcal{T}} b_{h}$, due to [AF], is known as the discrete Helmholtz decomposition. [In comparison with this, the version of [AF] interchanges the role of nonconforming and conforming terms; a change of the two components plus one change of signs proves the two versions equivalent.]
Definition 2.1. Given $p_{h}:=\nabla_{\mathcal{T}} u_{h}$ define the linear functional

$$
\mathcal{R e s}_{N C}: H^{1}(\Omega)^{m} \rightarrow \mathbb{R} \quad \text { by } \quad \mathcal{R e} s_{N C}(v):=\int_{\Omega} p_{h} \cdot \operatorname{Curl} v d x
$$

for $v \in H^{1}(\Omega)^{m}$ and for $m:=1$ if $n=2$ and $m:=3$ if $n=3$. Set

$$
\left\|\mathcal{R} e s_{N C}\right\|:=\sup _{v \in H^{1}(\Omega)^{m}, \operatorname{Curl} v \neq 0} \int_{\Omega} p_{h} \cdot \operatorname{Curl} v d x /\|\operatorname{Curl} v\|_{L} .
$$

The following result relates $\left\|\mathcal{R} e s_{N C}\right\|$ to (1.8)-(1.9). Notice carefully that $\tilde{u}_{h}$ is an arbitrary element in $H$ (and not necessarily some discrete function).
Theorem 2.2. There holds

$$
\min _{\bar{u}_{h} \in H}\left\|p_{h}-\nabla \tilde{u}_{h}\right\|_{L}=\left\|\mathcal{R} e s_{N C}\right\| \quad \text { and } \quad \mathcal{S}^{1}(\mathcal{T})^{m} \subset \operatorname{ker} \mathcal{R} e s_{N C} .
$$

Proof. The Helmholtz decomposition

$$
\begin{equation*}
p_{h}=\nabla a+\operatorname{Curl} b \tag{2.20}
\end{equation*}
$$

holds for a unique $a \in H_{0}^{1}(\Omega)=H$ and some $b \in H^{1}(\Omega)^{m}$ with the $L^{2}$ orthogonality of $\operatorname{Curl}\left(H^{1}(\Omega)^{m}\right)$ and $\nabla(H)$. Then

$$
\begin{align*}
\min _{\tilde{u}_{h} \in H}\left\|p_{h}-\nabla \tilde{u}_{h}\right\|_{L}^{2} & =\|\operatorname{Curl} b\|_{L}^{2} \\
& =\int_{\Omega} p_{h} \cdot \operatorname{Curl} b d x=\mathcal{R}^{2} s_{N C}(b)  \tag{2.21}\\
& \leq\left\|\mathcal{R} e s_{N C}\right\|\|\operatorname{Curl} b\|_{L}
\end{align*}
$$

For any $v \in H^{1}(\Omega)^{m}$ there holds, with (2.20) and $\int_{\Omega} \nabla a \cdot \operatorname{Curl} v d x=0$, that

$$
\begin{align*}
\operatorname{Res}_{N C}(v) & =\int_{\Omega}(\nabla a+\operatorname{Curl} b) \cdot \operatorname{Curl} v d x \\
& =\int_{\Omega} \operatorname{Curl} b \cdot \operatorname{Curl} v d x  \tag{2.22}\\
& \leq\|\operatorname{Curl} v\|_{L}\|\operatorname{Curl} b\|_{L} .
\end{align*}
$$

The combination of (2.21)-(2.22) shows

$$
\left\|\mathcal{R} e s_{N C}\right\|=\|\operatorname{Curl} b\|_{L}=\min _{\tilde{u}_{h} \in H}\left\|p_{h}-\nabla \tilde{u}_{h}\right\|_{L}
$$

This proves the first assertion of the theorem. The second is simpler in 2D and so solely shown for $n=m=3$ and $v_{h} \in \mathcal{S}^{1}(\mathcal{T})^{3}$. An elementwise integration by parts yields

$$
\begin{aligned}
\mathcal{R e s}_{N C}\left(v_{h}\right) & =\sum_{T \in \mathcal{T}} \int_{T} \nabla u_{h} \cdot \operatorname{Curl} v_{h} d x \\
& =\sum_{T \in \mathcal{T}} \int_{\partial T} u_{h}\left(\operatorname{Curl} v_{h}\right) \cdot \nu d s \\
& =\sum_{E \in \mathcal{E}} \int_{E}\left[u_{h} \operatorname{Curl} v_{h}\right] \cdot \nu_{E} d s
\end{aligned}
$$

for the unit normal $\nu_{E}$ on the element face $E$ and the jump [•] across $E$. Recall that $\int_{E}\left[u_{h}\right] d s=0$ by construction of the nonconforming finite element space $\mathcal{S}_{0}^{1, N C}(\mathcal{T})$. We claim that $\left[\operatorname{Curl} v_{h}\right] \cdot \nu_{E}=0$ on an interior face $E \in \mathcal{E}_{\Omega}:=\{E \in \mathcal{E}: E \not \subset \partial \Omega\}$. Since $v_{h}$ is a polynomial on $T_{+}$and $T_{-} \in \mathcal{T}, E=T_{+} \cap T_{-}$, and continuous along $E$, there holds ( $\otimes$ denotes the dyadic product)

$$
\left.D v_{h}\right|_{T_{+}}-\left.D v_{h}\right|_{T_{-}}=a \otimes \nu_{E} \quad \text { on } E
$$

for some polynomial $a$ in three components on $E$. A direct calculation shows that, therefore, the jump of $\left(\operatorname{Curl} v_{h}\right) \cdot \nu_{E}$ along $E$ vanishes. This proves our claim. We conclude continuity of $\left(\operatorname{Curl} v_{h}\right) \cdot \nu_{E}$ along $E \in \mathcal{E}$ and so

$$
\mathcal{R e} s_{N C}\left(v_{h}\right)=\sum_{E \in \mathcal{E}} \int_{E}\left[u_{h}\right]\left(\operatorname{Curl} v_{h}\right) \cdot \nu_{E} d s
$$

Since Curl $v_{h}$ is constant along $E$ and $\int_{E}\left[u_{h}\right] d s=0$ we conclude

$$
\mathcal{R e s}_{N C}\left(v_{h}\right)=0 .
$$

Remark 2.6. For $n=2$ dimensions, $|\operatorname{Curl} v|=|\nabla v|$ and $p_{h} \cdot \operatorname{Curl} v=$ $p_{h}^{\perp} \cdot \nabla v$ for $p_{h}^{\perp}:=\left(-p_{h 2}, p_{h 1}\right)$. Hence

$$
\begin{align*}
\mathcal{R} e s_{N C}(v) & =\int_{\Omega} p_{h}^{\perp} \cdot \nabla v d x \text { and } \\
\left\|\mathcal{R} e s_{N C}\right\| & =\sup _{v \in H^{1}(\Omega) \backslash \mathbb{R}} \int_{\Omega} p_{h}^{\perp} \cdot \nabla v d x /\|\nabla v\|_{L} \tag{2.23}
\end{align*}
$$

| Element | $M_{k}(T)$ | $D_{k}(T)$ |
| :---: | :---: | :---: |
| RT | $\mathcal{P}_{k}^{2}+x \cdot \mathcal{P}_{k}$ | $\mathcal{P}_{k}$ |
| BDM | $\mathcal{P}_{k+1}^{2}$ | $\mathcal{P}_{k}$ |
| BDFM | $\left\{q \in \mathcal{P}_{k+1}^{2}:\left.(q \cdot n)\right\|_{\partial T} \in \mathcal{R}_{k}(\partial T)\right\}$ | $\mathcal{P}_{k}$ |

Table 1. Standard 2D Mixed FEMs allowed in Theorem 2.4. Here, $\mathcal{P}_{k}$ denotes polynomials of total degree at most $k=0,1,2, \ldots$ and $\mathcal{R}_{k}(\partial T)$ denotes (not necessarily continuous) functions on $\partial T$ which equal a polynomial of degree at most $k$ on each edge.
is the usual operator norm as in (2.10). Notice the differences in the boundary conditions in (2.10) (where $v=0$ on $\partial \Omega$ ) and (2.23) (where $v$ has integral mean zero on $\Omega$ ). Since $\operatorname{div}_{\mathcal{T}} p_{h}^{\perp}=0$, we obtain

$$
\begin{equation*}
\eta_{R}^{(N C)}:=\left\|h_{\mathcal{E}}^{1 / 2}\left[p_{h} \cdot \tau_{\mathcal{E}}\right]\right\|_{L^{2}(\cup \mathcal{E})} \tag{2.24}
\end{equation*}
$$

(with the piecewise tangential unit vector $\tau_{\mathcal{E}}$ ) instead of (2.11)-(2.12) and, as in (2.13),

$$
\begin{equation*}
\eta_{A}:=\left\|p_{h}-\sum_{z \in \mathcal{N}}\left(A_{h} p_{h}\right) \varphi_{z}\right\|_{L^{2}(\Omega)} . \tag{2.25}
\end{equation*}
$$

Remark 2.7. We stress that all other estimators, for instance the localized or equilibrated implicit estimators of [AO, BS], are available for the assessment of

$$
\begin{equation*}
\left\|\mathcal{R} e s_{N C}\right\| \approx \eta_{R}^{(N C)} \lesssim \eta_{A} \tag{2.26}
\end{equation*}
$$

as well. The averaging estimator (2.25) concerns discontinuities in normal and tangential components and so

$$
\begin{equation*}
\left\|\mathcal{R e}_{N C}\right\|+\left\|\mathcal{R} e s_{H}\right\|_{H}+\left\|h_{\mathcal{T}}^{2} \nabla g\right\|_{L^{2}(\Omega)} \approx \eta_{A}+\left\|h_{\mathcal{T}}^{2} \nabla_{\mathcal{T}} g\right\|_{L^{2}(\Omega)} \tag{2.27}
\end{equation*}
$$

Remark 2.8. The situation for $n=3$ dimensions is more delicate and we refer to $[\mathrm{CBJ}]$ for reliable and efficient explicit error estimators.
2.4. Mixed Finite Element Methods. The Laplace equation is split into $\operatorname{div} p+g=0$ and the weak form of $p=\nabla u$. The resulting mixed formulation involves a bilinear form as in (1.5). Its discrete version involves

$$
p_{h} \in L_{h} \subseteq \mathcal{L}^{k+1}(\mathcal{T})^{n} \cap H(\operatorname{div} ; \Omega) \quad \text { and } \quad u_{h} \in H_{h} \subseteq \mathcal{L}^{k}(\mathcal{T})
$$

for $k=0,1$. The couple ( $p_{h}, u_{h}$ ) is supposed to satisfy

$$
\begin{align*}
& \int_{\Omega} p_{h} \cdot q_{h} d x+\int_{\Omega} u_{h} \operatorname{div} q_{h} d x=0 \quad \text { for all } q_{h} \in L_{h} \\
& \int_{\Omega} v_{h} \operatorname{div} p_{h} d x=-\int_{\Omega} g v_{h} d x \quad \text { for all } v_{h} \in \mathcal{L}^{k}(\mathcal{T}) . \tag{2.28}
\end{align*}
$$

With the above sets $D_{k}(T)$ and $M_{k}(T)$ from Table 1 for $n=2$ we define

$$
\begin{aligned}
& L_{h}:=M_{k}(\mathcal{T}):=\left\{q_{h} \in H(\operatorname{div} ; \Omega): \forall T \in \mathcal{T},\left.q_{h}\right|_{T} \in M_{k}(T)\right\}, \\
& H_{h}:=D_{k}(\mathcal{T}):=\left\{v_{h} \in L^{\infty}(\Omega): \forall T \in \mathcal{T},\left.v_{h}\right|_{T} \in D_{k}(T)\right\} .
\end{aligned}
$$

Theorem 2.1 is applied to estimate $\left\|p-p_{h}\right\|_{L}+\left\|u-\tilde{u}_{h}\right\|_{H}$ for some $\tilde{u}_{h} \in H$ different from $u_{h}$. The evaluation of $\mathcal{R} e s_{H}$ follows the arguments of the conforming finite element situation in Subsection 2.2.

Theorem 2.3. Given $g \in H^{1+k}(\mathcal{T}):=\left\{g \in L^{2}(\Omega): \forall T \in \mathcal{T},\left.g\right|_{T} \in\right.$ $\left.H^{1+k}(T)\right\}$ and $k=0,1$,

$$
\left\|\mathcal{R} \operatorname{es}_{H}\right\|_{H^{*}} \lesssim\left\|h_{\mathcal{T}}^{2+k} \nabla_{\mathcal{T}}^{1+k} g\right\|_{L^{2}(\Omega)} \quad \text { is of higher order. }
$$

Proof. An integration by parts shows, for all $v \in H$,

$$
\mathcal{R} e s_{H}(v):=-\int_{\Omega} g v d x+\int_{\Omega} p_{h} \cdot \nabla v d x=-\int_{\Omega} v\left(g+\operatorname{div} p_{h}\right) d x .
$$

Note that there are no jump terms across interior element boundaries since $p_{h} \in H(\operatorname{div} ; \Omega)$ (and, equivalently, $\left.\left[p_{h}\right] \cdot \nu_{\mathcal{E}}=0\right)$. In the lowestorder cases, $H_{h}=\mathcal{L}^{0}(\mathcal{T})$ and $(2.28)_{b}$ lead to

$$
-\left.\operatorname{div} p_{h}\right|_{T}=\int_{T} g d x /|T|=:\left.g_{h}\right|_{T} \quad \text { for all } T \in \mathcal{T}
$$

Consequently, if $v_{h}$ and $g_{h}$ denote the $\mathcal{T}$-piecewise constant averages of $v$ and $g$, respectively, Poincaré inequalities show

$$
\begin{align*}
\mathcal{R e s}_{H}(v) & =-\int_{\Omega}\left(v-v_{h}\right)\left(g-g_{h}\right) d x \\
& \leq\left\|h_{\mathcal{T}}\left(g-g_{h}\right)\right\|_{L^{2}(\Omega)}\left\|\left(v-v_{h}\right) / h_{\mathcal{T}}\right\|_{L^{2}(\Omega)}  \tag{2.29}\\
& \lesssim\left\|h_{\mathcal{T}}^{2} \nabla g\right\|_{L^{2}(\Omega)}\|\nabla v\|_{L^{2}(\Omega)} .
\end{align*}
$$

The proof is finished for $k=0$. For $k=1$, the second equation in (2.28) implies on each element domain $T \in \mathcal{T}$ that $-\operatorname{div} p_{h}=g_{h}:=\Pi_{1} g$ is the $L^{2}$ projection of $g$ onto $P_{1}(T)$. Hence (2.29) can be moified to yield the upper bound $\left\|h_{\mathcal{T}}^{3} \nabla_{\mathcal{T}}^{2} g\right\|_{L^{2}(\Omega)}\|\nabla v\|_{L^{2}(\Omega)}$.

For the evaluation of $\left\|\mathcal{R} e s_{L}\right\|_{L^{*}}$ we have the following analogue of Theorem 2.2 which shows (1.8)-(1.9) in the notation from Definition 2.1;
recall that $\tilde{u}_{h}$ denotes an arbitrary (not necessarily discrete) element in $H$.

Theorem 2.4. For each of the mixed FEM of Table 1 there holds

$$
\min _{\bar{u}_{h} \in H}\left\|p_{h}-\nabla \tilde{u}_{h}\right\|_{L}=\left\|\mathcal{R} e s_{N C}\right\| \quad \text { and } \quad \mathcal{S}^{1}(\mathcal{T})^{m} \subset \operatorname{ker} \mathcal{R} e s_{N C} .
$$

Proof. The assertion follows as in Theorem 2.2: Given $v_{h} \in \mathcal{S}^{1}(\mathcal{T})^{m}$ the second part of its proof showed Curl $v_{h} \in H(\operatorname{div} ; \Omega)$ by $\left[\operatorname{Curl} v_{h}\right] \cdot \nu_{E}=0$. Hence,

$$
q_{h}:=\operatorname{Curl} v_{h} \in \mathcal{L}^{1}(\mathcal{T})^{n} \cap H(\operatorname{div}, \Omega) \subset L_{h} .
$$

The last inclusion holds for the finite element spaces of Table 1. Thus, $q_{h}:=\operatorname{Curl} v_{h}$ may be considered in $(2.28)_{a}$ and shows

$$
\mathcal{R} e s_{N C}\left(v_{h}\right)=\int_{\Omega} p_{h} \cdot \operatorname{Curl} v_{h} d x=-\int_{\Omega} u_{h} \operatorname{div} \operatorname{Curl} v_{h} d x=0
$$

Remark 2.9. Based on Theorem 2.4, the evaluation of $\left\|\mathcal{R} e s_{N C}\right\|$ follows the arguments of Subsections 2.2 and 2.3; e.g. for $n=2, \eta_{R}^{(N C)}$, and $\eta_{A}$ from (2.24)-(2.25) there holds

$$
\left\|\mathcal{R} e s_{N C}\right\| \approx \eta_{R}^{(N C)} \approx \eta_{A} .
$$

## 3. Applications to the Stokes Problem

The stationary incompressible fluid flow can be modelled by the Stokes equations: Given $g \in L^{2}(\Omega)$ seek $(u, p) \in H \times L:=H_{0}^{1}(\Omega)^{n} \times$ $L_{0}^{2}(\Omega)$ with

$$
\begin{align*}
& \int_{\Omega} 2 \mu \varepsilon(u): \varepsilon(v) d x-\int_{\Omega} p \operatorname{div} v d x=\int_{\Omega} g v d x \\
& -\int_{\Omega} q \operatorname{div} u d x=0 \quad \text { for all }(v, q) \in H_{0}^{1}(\Omega)^{n} \times L_{0}^{2}(\Omega) \tag{3.1}
\end{align*}
$$

Here, $L_{0}^{2}(\Omega):=\left\{q \in L^{2}(\Omega): \int_{\Omega} q d x=0\right\} \equiv L^{2}(\Omega) / \mathbb{R}$ fixes a global additive constant in the pressure (because of lacking Neumann boundary conditions). In this case, (3.1) is equivalent to a formulation with the non-symmetric gradient $\nabla u$ instead of its symmetric part

$$
\begin{equation*}
\varepsilon(u):=\left(\nabla u+\nabla u^{T}\right) / 2 \in L:=L^{2}\left(\Omega ; \mathbb{R}_{\mathrm{sym}}^{n \times n}\right) \tag{3.2}
\end{equation*}
$$

here; $\mathbb{R}_{\text {sym }}^{n \times n}=\left\{A \in \mathbb{R}^{n \times n}: A=A^{T}\right\}$. (Colon denotes the Euclidean scalar product, $A: B=\sum_{j, k=1}^{n} A_{j k} B_{j k}$ for $A, B \in \mathbb{R}^{n \times n}$.)

It is well-known that (3.1) has a unique solution $(u, p)$. We discuss conforming and nonconforming finite element approximations of the

Stokes equations. Given a regular triangulation $\mathcal{T}$ and a $\mathcal{T}$-piecewise $H^{1}$ function (written $v \in H^{1}(\mathcal{T})$ ) set

$$
\left.\varepsilon_{\mathcal{T}}(v):=\left(\nabla_{\mathcal{T}} v+\nabla_{\mathcal{T}} v^{T}\right)\right) / 2 \in L:=L^{2}\left(\Omega ; \mathbb{R}_{\mathrm{sym}}^{n \times n}\right)
$$

for the $\mathcal{T}$-piecewise gradient $\nabla_{\mathcal{T}} ; \nabla_{\mathcal{T} v}$ equals $\nabla\left(\left.v\right|_{T}\right)$ on each $T \in$ $\mathcal{T}$; let $\operatorname{div}_{\mathcal{T}}$ denote the $\mathcal{T}$-piecewise divergence operator. To describe conforming and nonconforming finite element methods simultaneously, suppose $u_{h} \in \mathcal{L}^{k}(\mathcal{T})^{n}$ and $p_{h} \in \mathcal{L}^{k}(\mathcal{T}) \cap L_{0}^{2}(\Omega)$ satisfy

$$
\begin{equation*}
\int_{\Omega} 2 \mu \varepsilon_{\mathcal{T}}\left(u_{h}\right): \varepsilon_{\mathcal{T}}\left(v_{h}\right) d x-\int_{\Omega} p_{h} \operatorname{div}_{\mathcal{T}} v_{h} d x=\int_{\Omega} g v_{h} d x \tag{3.3}
\end{equation*}
$$

for all $v_{h} \in \mathcal{S}_{0}^{1}(\mathcal{T})^{n}$.
Remark 3.1. Even for nonconforming schemes we suppose that (3.3) holds for a continuous test function $v_{h}$. For the lowest-order finite element schemes, this implies the restriction to triangular finite elements.

Remark 3.2. Throughout the discussion of this paper, the discrete $u_{h}$ and $p_{h}$ are supposed to be piecewise polynomials of some degree $\leq k$. This does not mean that we propose some $P_{k}^{n} \times P_{k}$ finite element method - they may be instable. However, our a posteriori analysis partly includes error control even for unstable methods.

Remark 3.3. The condition $(3.1)_{b}$ has no discrete analog in (3.3) because that is not needed in our a posteriori error analysis. However, since $\left\|\operatorname{div}_{\mathcal{T}} u_{h}\right\|_{L^{2}(\Omega)}$ arises in estimates of $\mathcal{R} e s_{L}$ below, it is understood below that $\left\|\operatorname{div}_{\mathcal{T}} u_{h}\right\|_{L^{2}(\Omega)}$ is small.
Remark 3.4. The list of examples for $n=2$ includes conforming finite elements such as the MINI element, the $P_{2}-P_{0}$ finite element, and the Taylor-Hood element [BF] and the nonconforming finite element due to Kouhia and Stenberg [KS].

Remark 3.5. There are also finite element methods for the unsymmetric formulation [BF] such as the popular Crouzeix-Raviart finite element. Since the arguments of this section work verbatim (if not simpler), we omit details and refer to [DDP, V1].
Definition 3.1. Given $u_{h} \in H^{1}(\mathcal{T})^{n}$ and $p_{h} \in L_{0}^{2}(\Omega)$ set

$$
\begin{equation*}
\sigma_{h}:=2 \mu \varepsilon_{\mathcal{T}}\left(u_{h}\right)-p_{h} \mathbf{1} \in L:=L^{2}\left(\Omega ; \mathbb{R}_{\mathrm{sym}}^{n \times n}\right) \tag{3.4}
\end{equation*}
$$

and define the linear functional $\mathcal{R e s}_{H}: H \rightarrow \mathbb{R}$ by

$$
\begin{equation*}
\mathcal{R} e s_{H}(v)=\int_{\Omega}\left(g \cdot v-\sigma_{h}: \varepsilon(v)\right) d x \quad \text { for } v \in H:=H_{0}^{1}(\Omega)^{n} \tag{3.5}
\end{equation*}
$$

The linear function space $H$ is endowed with the norm $\|v\|_{H}:=|v|_{1,2}$ $:=\|\nabla v\|_{2}$ for $v \in H$ such that

$$
\left\|\mathcal{R} e s_{H}\right\|_{H^{*}}=\sup _{v \in H \backslash\{0\}} \mathcal{R} e s_{H}(v) /|v|_{1,2}
$$

The residual $\left\|\mathcal{R} e s_{L}\right\|_{L^{*}}$ involves the deviatoric-part operator

$$
\begin{equation*}
\operatorname{dev} A=A-(\operatorname{tr}(A) / n) \mathbf{1} \quad \text { for any } A \in \mathbb{R}^{n \times n} \tag{3.6}
\end{equation*}
$$

(where $\operatorname{tr}(A)=A_{11}+\cdots+A_{n n}$ is the trace of $A$ ).
Theorem 3.1. Let $(u, p) \in H \times L_{0}^{2}(\Omega)$ solve (3.1) and set $\sigma:=$ $2 \mu \varepsilon(u)-p \mathbf{1} \in L$. Let $\left(u_{h}, p_{h}\right)$ satisfy (3.3) and define $\sigma_{h}$ as in (3.4). Then, for any $\tilde{u}_{h} \in H$, there holds

$$
\begin{equation*}
\left\|\sigma-\sigma_{h}\right\|_{L}+\left\|u-\tilde{u}_{h}\right\|_{H} \approx\left\|\varepsilon\left(\tilde{u}_{h}\right)-\operatorname{dev} \varepsilon_{\mathcal{T}}\left(u_{h}\right)\right\|_{L}+\left\|\operatorname{Re}_{H}(v)\right\|_{H^{*}} . \tag{3.7}
\end{equation*}
$$

Before we focus on its proof, we briefly comment on applications of the theorem. The residual $\mathcal{R} e s_{H}$ satisfies $\mathcal{R} e s_{H}\left(v_{h}\right)=0$ for all $v_{h} \in \mathcal{S}_{0}^{1}(\mathcal{T})^{n}$ and, since $\sigma_{h}$ is symmetric, can be recast into

$$
\begin{equation*}
\mathcal{R e s}_{H}(v)=\int_{\Omega}\left(g \cdot v-\sigma_{h}: \nabla v\right) d x \tag{3.8}
\end{equation*}
$$

That is, $\mathcal{R e s}_{H}$ in (3.5) is the sum of $j=1,2, \ldots, n$ residuals $\mathcal{R} e s\left(v e_{j}\right)$ of the form in $(2.6)_{b}$ where $e_{j}$ is the $j$-th canonical unit vector in $\mathbb{R}^{n}$ and, here, $v$ in $H_{0}^{1}(\Omega)$ is a scalar. As a consequence, the residual evaluation can follow the same lines as in Subsection 2.1.

The discussion of $\left\|\mathcal{R} e s_{L}\right\|_{L^{*}}=\left\|\varepsilon\left(\tilde{u}_{h}\right)-\operatorname{dev} \varepsilon_{\mathcal{T}}\left(u_{h}\right)\right\|_{L}$ follows two cases. In case I, for any conforming approximation $u_{h} \in H$, the choice $\tilde{u}_{h}=u_{h}$ yields

$$
\begin{equation*}
\left\|\mathcal{R} e s_{L}\right\|_{L^{*}}=\left\|\varepsilon\left(u_{h}\right)-\operatorname{dev} \varepsilon_{\mathcal{T}}\left(u_{h}\right)\right\|_{L}=n^{-1}\left\|\operatorname{div} u_{h}\right\|_{L^{2}(\Omega)} \tag{3.9}
\end{equation*}
$$

This is an appropriate error contribution and, at the same time, an error estimator.

In case II, $u_{h} \notin H$ and the estimation of the nonconformity terms is analogous to that of Subsection 2.3 but slightly more involved because of the interaction of the divergence residual and the $u_{h}-\tilde{u}_{h}$ approximation. If one accepts $\left\|\operatorname{div}_{\mathcal{T}} u_{h}\right\|_{L^{2}(\Omega)}$ as a proper error term (cf. Remark 3.3), the upper bound

$$
\begin{equation*}
\left\|\mathcal{R} e s_{L}\right\|_{L^{*}} \leq\left\|\varepsilon_{\mathcal{T}}\left(u_{h}-\tilde{u}_{h}\right)\right\|_{L}+n^{-1 / 2}\left\|\operatorname{div}_{\mathcal{T}} u_{h}\right\|_{L^{2}(\Omega)} \tag{3.10}
\end{equation*}
$$

can be minimised according to a symmetric form of a Helmholtz decomposition. To quote results from the literature let $n=2$ for a moment.

Lemma 3.2. ([CD, Lemma 3.2]). Given any $\tau \in L^{2}\left(\Omega ; \mathbb{R}_{\mathrm{sym}}^{2 \times 2}\right)$ there exist $a \in H_{0}^{1}(\Omega)^{2}$ and $b \in H^{2}(\Omega)$ such that

$$
\begin{equation*}
\tau=\varepsilon(a)+\operatorname{Curl} \operatorname{Curl} b \tag{3.11}
\end{equation*}
$$

(where Curl Curl $b$ has the four entries $b_{, 22},-b_{, 12},-b_{, 12}, b_{, 11}$ where $b_{, \alpha \beta}=$ $\left.\partial^{2} b / \partial x_{\alpha} \partial x_{\beta}\right)$.

Suppose there holds $\tau=\varepsilon_{\mathcal{T}}\left(u_{h}\right)=\varepsilon(a)+$ Curl Curl $b$ in (3.11). The, the lemma suggests the choice $\tilde{u}_{h}=a$ and, as for Theorem 2.2, one proves

$$
\begin{equation*}
\min _{\tilde{u}_{h} \in H}\left\|\varepsilon\left(u_{h}-\tilde{u}_{h}\right)\right\|_{L}=\|\operatorname{Curl} \operatorname{Curl} b\|_{L}=\left\|\mathcal{R} e s_{N C}\right\|_{H^{*}} \tag{3.12}
\end{equation*}
$$

for

$$
\mathcal{R} e s_{N C}(v):=\int_{\Omega} \varepsilon_{\mathcal{T}}\left(u_{h}\right): \operatorname{Curl} v d x \text { for all } v \in H
$$

[The proof follows closely (2.20)-(2.22) and is hence omitted.] More details may be found in [CF1] where it is in particular shown (for the Kouhia-Stenberg FEM and $n=2$ ) that

$$
\begin{equation*}
\left\|\mathcal{R} e s_{N C}\right\|_{H^{*}} \lesssim\left\|h_{\mathcal{E}}^{1 / 2}\left[\nabla u_{h} \cdot \tau_{\mathcal{E}}\right]\right\|_{L^{2}(\cup \mathcal{E})} . \tag{3.13}
\end{equation*}
$$

The remaining part of this section is devoted to the proof of Theorem 3.1 for $n=2$ or $n=3$. To employ the mixed approach of Subsection 1.3 set

$$
\begin{align*}
& a(\sigma, \tau):=\int_{\Omega} 1 /(2 \mu) \operatorname{dev} \sigma: \operatorname{dev} \tau d x \quad \text { for } \sigma, \tau \in L  \tag{3.14}\\
& b(\sigma ; v):=-\int_{\Omega} \sigma: \varepsilon(v) d x \quad \text { for }(\sigma, v) \in L \times H .
\end{align*}
$$

Let $L / \mathbb{R}:=\left\{\sigma \in L: \int_{\Omega} \operatorname{tr} \sigma d x=0\right\}$.
Lemma 3.3. The operator $A: X \rightarrow X^{*}$, defined for $(\sigma, u) \in X:=$ $(L / \mathbb{R}) \times H$ by

$$
\begin{equation*}
(A(\sigma, u))(\tau, v):=a(\sigma, \tau)+b(\sigma, v)+b(\tau, u) \tag{3.15}
\end{equation*}
$$

is linear, bounded and bijective. [This result holds for $n=2,3$.]
Proof. The bijectivity of $A$ is the only not so immediate part of the lemma. Proposition 3.1 in [BF, Chapter IV] states

$$
\|\sigma\|_{L} \lesssim\|\operatorname{dev} \sigma\|_{L}+\|\operatorname{div} \sigma\|_{L^{2}(\Omega)} \quad \text { for all } \sigma \in L / \mathbb{R}
$$

Since any $\sigma \in L / \mathbb{R}$ with $b(\sigma ; \cdot)=0($ written $\sigma \in \operatorname{ker} B)$ satisfies $\operatorname{div} \sigma=$ 0 this implies

$$
\|\sigma\|_{L}^{2} \lesssim a(\sigma, \sigma) \quad \text { for all } \sigma \in \operatorname{ker} B
$$

Hence $a$ is elliptic on ker $B$. This is one of the main ingredients of the general theory on mixed finite element mehods $[\mathrm{BF}]$. The remaining details are omitted.

Proof of Theorem 3.1. The inf-sup condition for $A$ follows from Lemma 3.3. The resulting equivalence (1.7) reads

$$
\left\|\sigma-\sigma_{h}\right\|_{L}+\left\|u-\tilde{u}_{h}\right\|_{H} \approx\left\|\mathcal{R} e s_{L}\right\|_{L^{*}}+\left\|\operatorname{Res}_{H}\right\|_{H^{*}}
$$

The residuals on the right-hand side result from (1.6). In particular, if we employ $\operatorname{dev} \varepsilon(u)=\varepsilon(u)$ (from $\operatorname{div} u=0$ ), there holds

$$
\begin{aligned}
\mathcal{R} e s_{L}(\tau) & =\int_{\Omega}\left(\left(1 /(2 \mu) \operatorname{dev}\left(\sigma-\sigma_{h}\right)-\varepsilon\left(u-\tilde{u}_{h}\right)\right): \tau d x\right. \\
& =\int_{\Omega}\left(\varepsilon\left(\tilde{u}_{h}\right)-\operatorname{dev}_{\mathcal{T}} \varepsilon\left(u_{h}\right)\right): \tau d x
\end{aligned}
$$

The remaining details are omitted.

## 4. Applications to linear elasticity

This section is devoted to the Navier-Lamè equations and its conforming, nonconforming, and mixed finite element approximation. An analysis of enhanced finite elements in the same framework is given in [BCR]. It is an important feature of the presented unifying theory that the resulting a posteriori error estimates are robust with respect to the Lamè parameter $\lambda \rightarrow \infty$.
4.1. Model Problem. We adopt the notation of the previous two sections and continue with a linear stress-strain relation of the form

$$
\mathbb{C} A:=\lambda \operatorname{tr}(A) \mathbf{1}+2 \mu A \quad \text { for } A \in \mathbb{R}^{n \times n}
$$

with inverse relation

$$
\mathbb{C}^{-1} A=1 /(2 \mu) A-\lambda /(2 \mu(n \lambda-2 \mu)) \operatorname{tr}(A) \mathbf{1} \quad \text { for } A \in \mathbb{R}^{n \times n}
$$

In the continuous model, $\sigma=\mathbb{C} \varepsilon(u) \in L:=L^{2}\left(\Omega ; \mathbb{R}_{\text {sym }}^{n \times n}\right)$ and the resulting model problem reads: Given $g \in L^{2}(\Omega)^{n}$ find $u \in H:=$ $H_{0}^{1}(\Omega)^{n}$ with

$$
\begin{equation*}
g+\operatorname{div} \mathbb{C} \varepsilon(u)=0 \quad \text { in } \Omega \tag{4.1}
\end{equation*}
$$

The material parameters $\lambda$ and $\mu$ are positive and hence (4.1) is an elliptic PDE with a unique solution $u$.

To employ the unified theory of Subsection 1.3 let

$$
\begin{equation*}
a(\sigma, \tau):=\int_{\Omega}\left(\mathbb{C}^{-1} \sigma\right): \tau d x \quad \text { for } \sigma, \tau \in L:=L^{2}\left(\Omega ; \mathbb{R}_{\mathrm{sym}}^{n \times n}\right) \tag{4.2}
\end{equation*}
$$

replace $(3.14)_{a}$ and adopt $b(\sigma ; v)$ from $(3.14)_{b}$.
Lemma 4.1 ([BCR]). The operator $A: X \rightarrow X^{*}$ from (3.15) is linear, bounded, and bijective and the operator norms of $A$ and $A^{-1}$ are $\lambda$ independent.

Remark 4.1. The operator $A$ represents the weak form of the HellingerReissner principle in mechanics. Lemma 4.1 is the analogue of Lemma 3.3.
4.2. Conforming Finite Element Methods. Although the (lower order) conforming finite element methods are not robust in $\lambda \rightarrow \infty$ we introduce a robust error estimation. Given a finite element approximation $u_{h} \in H=H_{0}^{1}(\Omega)^{n}$ with

$$
\begin{equation*}
\int_{\Omega} \varepsilon\left(u_{h}\right): \mathbb{C} \varepsilon\left(v_{h}\right) d x=\int_{\Omega} g \cdot v_{h} d x \quad \text { for all } v_{h} \in \mathcal{S}_{0}^{1}(\mathcal{T})^{n} \tag{4.3}
\end{equation*}
$$

let $e:=u-u_{h} \in H_{0}^{1}(\Omega)$ (where $u$ solves (4.1)) denote the error.
Theorem 4.2. With $\lambda$-independent constants in $\approx$, there holds
$\|\mathbb{C} \varepsilon(e)\|_{L^{2}(\Omega)}+\|e\|_{H^{1}(\Omega)} \approx \sup _{v \in H \backslash\{0\}} \int_{\Omega}\left(g \cdot v-\varepsilon\left(u_{h}\right): \mathbb{C} \varepsilon(v)\right) d x /\|v\|_{H^{1}(\Omega)}$.
Proof. With $\tilde{u}_{h}=u_{h} \in H, \sigma:=\mathbb{C} \varepsilon(u)$, and $\sigma_{h}:=\mathbb{C} \varepsilon\left(u_{h}\right)$, Equivalence (1.7) results in

$$
\left\|\sigma-\sigma_{h}\right\|_{L}+\|e\|_{H^{1}(\Omega)} \approx\left\|\mathcal{R} e s_{H}\right\|_{H^{*}}
$$

where $\operatorname{Res}_{H}(v)$ is defined in (3.8). This implies the assertion.

As in the previous section (cf. the discussion about (3.8)), the estimation of $\left\|\mathcal{R} e s_{H}\right\|_{H^{*}}$ follows the lines of Subsection 2.1. Given any estimator $\eta$ (of the various choices (2.11)-(2.14)) with

$$
\begin{equation*}
\left\|\mathcal{R} e s_{H}\right\|_{H^{*}} \approx \eta \quad \text { (up to h.o.t.) } \tag{4.4}
\end{equation*}
$$

the estimate [the proof follows from (4.4) and $\left.\|e\|_{H^{1}(\Omega)} \lesssim\left\|\sigma-\sigma_{h}\right\|_{L^{2}(\Omega)}\right]$

$$
\begin{equation*}
\left\|\sigma-\sigma_{h}\right\|_{L^{2}(\Omega)} \approx \eta \quad \text { (up to h.o.t.) } \tag{4.5}
\end{equation*}
$$

appears to be new (where h.o.t. refers to $\left\|h_{\tau}^{2} \nabla g\right\|_{L^{2}(\Omega)}$ for the first-order finite element schemes). The point is that the constants behind $\approx$ in (4.5) are $\lambda$-independent. This is different for the standard estimate

$$
\begin{equation*}
\left\|\mathbb{C}^{1 / 2} \varepsilon(e)\right\|_{L^{2}(\Omega)} \lesssim \eta \lesssim\left\|\sigma-\sigma_{h}\right\|_{L^{2}(\Omega)} \quad \text { (up to h.o.t.). } \tag{4.6}
\end{equation*}
$$

In fact, for the proof of (4.6) one argues

$$
\begin{aligned}
& \int_{\Omega} \varepsilon(e): \mathbb{C} \varepsilon(e) d x=\operatorname{Res}_{H}(e) \leq\left\|\mathcal{R} e s_{H}\right\|_{H^{*}}\|e\|_{H} \\
& \quad \lesssim\left\|\mathbb{C}^{-1 / 2}\right\|_{\infty}\left\|\mathbb{C}^{1 / 2} \varepsilon(e)\right\|_{L^{2}(\Omega)} \eta \lesssim\left\|\mathbb{C}^{1 / 2} \varepsilon(e)\right\|_{L^{2}(\Omega)} \eta .
\end{aligned}
$$

Finally, inverse estimates verify $(4.6)_{b}$ up to higher order terms (h.o.t.). Notice that (4.5) is balanced in $\lambda$ while (4.6) is not. Thus (4.4) establishes $\lambda$-robust a posteriori error control of the $L^{2}$-stress error $\| \sigma-$ $\sigma_{h} \|_{L^{2}(\Omega)}$. A corresponding result for the energy norm $\left\|\mathbb{C}^{1 / 2} \varepsilon(e)\right\|_{L^{2}(\Omega)}$ remains as an open problem.
4.3. Nonconforming Finite Element Methods. Robust a priori convergence estimates are known for nonconforming finite element methods such as

$$
\mathcal{S}_{0}^{1}(\mathcal{T})^{2} \subset \mathcal{S}_{0}^{1}(\mathcal{T}) \times \mathcal{S}_{0}^{1, N C}(\mathcal{T}) \subset H_{0}^{1}(\Omega) \times H^{1}(\mathcal{T})
$$

due to Kouhia-Stenberg [KS]. Here and throughout this subsection, let $n=2$ and consider merely triangles. Suppose that the discrete solution $u_{h} \in H^{1}(\mathcal{T})^{2}$ satisfies

$$
\begin{equation*}
\int_{\Omega} \varepsilon_{\mathcal{T}}\left(u_{h}\right): \mathbb{C} \varepsilon\left(v_{h}\right) d x=\int_{\Omega} g \cdot v_{h} d x \quad \text { for all } v_{h} \in \mathcal{S}_{0}^{1}(\mathcal{T})^{2} \tag{4.7}
\end{equation*}
$$

and set $\sigma_{h}:=\mathbb{C} \varepsilon_{\mathcal{T}}\left(u_{h}\right), \sigma:=\mathbb{C} \varepsilon(u) \in L:=L^{2}\left(\Omega ; \mathbb{R}_{\mathrm{sym}}^{2 \times 2}\right)$.
Theorem 4.3. For any $\tilde{u}_{h} \in H$ there holds

$$
\begin{equation*}
\left\|\sigma-\sigma_{h}\right\|_{L}+\left\|u-\tilde{u}_{h}\right\|_{H} \approx\left\|\varepsilon_{\mathcal{T}}\left(u_{h}-\tilde{u}_{h}\right)\right\|_{L}+\left\|\mathcal{R e}^{2}\right\|_{H^{*}} \tag{4.8}
\end{equation*}
$$

Proof. This is a result of Equivalence (1.7) and Lemma 4.1. The details are similar to those of the previous sections and hence omitted.

The discussion of $\left\|\varepsilon_{\mathcal{T}}\left(u_{h}-\tilde{u}_{h}\right)\right\|_{L}$ follows the lines of (3.10)-(3.13) in the previous section and hence are omitted; cf. [CF2].
4.4. Mixed Finite Element Methods. Another $\lambda$-robust approximation is feasible with a mixed finite element method with reduced symmetry [ABD, S] for

$$
\sigma_{h} \in H(\operatorname{div} ; \Omega)
$$

i.e. $\operatorname{div} \sigma_{h} \in L^{2}\left(\Omega ; \mathbb{R}^{n}\right)$ and $\sigma_{h} \in \mathcal{L}^{k}\left(\mathcal{T} ; \mathbb{R}^{n \times n}\right)$ with, in general, $\operatorname{As}\left(\sigma_{h}\right)$ $:=\sigma_{h}-\operatorname{sym} \sigma_{h} \not \equiv 0, \operatorname{sym} \sigma_{h}:=\left(\sigma_{h}^{T}+\sigma_{h}\right) / 2$. Suppose that $\sigma_{h}$ satisfies

$$
\begin{equation*}
-\int_{\Omega} v_{h} \cdot \operatorname{div} \sigma_{h} d x=\int_{\Omega} v_{h} \cdot g d x \quad \text { for all } v_{h} \in \mathcal{L}^{0}(\mathcal{T})^{n} \tag{4.9}
\end{equation*}
$$

Theorem 4.4. For any $\tilde{u}_{h} \in H$ there holds

$$
\begin{align*}
& \left\|\sigma-\operatorname{sym} \sigma_{h}\right\|_{L}+\left\|u-\tilde{u}_{h}\right\|_{H}  \tag{4.10}\\
& \quad \approx\left\|\varepsilon\left(\tilde{u}_{h}\right)-\mathbb{C}^{-1} \operatorname{sym} \sigma_{h}\right\|_{L}+\left\|\mathcal{R e s} s_{H}\right\|_{H^{*}} .
\end{align*}
$$

Proof. Lemma 4.1 and Equivalence (1.7) apply to $\sigma-\operatorname{sym} \sigma_{h}$ and $u-\tilde{u}_{h}$. We omit the details.

Notice that the residual satisfies

$$
\begin{aligned}
\operatorname{Res}_{H}(v) & =\int_{\Omega}\left(v \cdot g-\varepsilon(v): \sigma_{h}\right) d x \\
& =\int_{\Omega} v \cdot\left(g+\operatorname{div} \sigma_{h}\right) d x-\int_{\Omega} A s\left(\sigma_{h}\right): \nabla(v) d x \quad \text { for } v \in H
\end{aligned}
$$

(For a proof observe that $\varepsilon(v): \operatorname{sym} \sigma_{h}=\nabla(v):\left(\sigma_{h}-A s\left(\sigma_{h}\right)\right)$ and employ an integration by parts.) The estimation of $\left\|\mathcal{R} e s_{H}\right\|_{H^{*}}$ may hence follow as in (2.29) and yields

$$
\begin{equation*}
\left\|\mathcal{R} e s_{H}\right\|_{H^{*}} \lesssim\left\|h_{\mathcal{T}}^{2} \nabla_{\mathcal{T}} g\right\|_{L^{2}(\Omega)}+\left\|A s\left(\sigma_{h}\right)\right\|_{L^{2}(\Omega)} \tag{4.11}
\end{equation*}
$$

The first term on the right-hand side of (4.11) is of higher order for $g \in H^{1}(\mathcal{T})^{n}$ and first order schemes such as PEERS. The estimation of

$$
\begin{equation*}
\min _{\tilde{u}_{h} \in H}\left\|\varepsilon\left(\tilde{u}_{h}\right)-\mathbb{C}^{-1} \operatorname{sym} \sigma_{h}\right\|_{L}=\left\|\mathcal{R} e s_{N C}\right\| \tag{4.12}
\end{equation*}
$$

follows closely the discussion of (3.10)-(3.13). We refer to [CD, CDFH] for the remaining details on the approximation of

$$
\mathcal{R e} s_{N C}(v):=\int_{\Omega} \operatorname{Curl} v: \operatorname{sym} \mathbb{C}^{-1} \sigma_{h} d x
$$

Notice that (4.10) concerns the symmetric part of the error $\left\|\sigma-\sigma_{h}\right\|_{L}$. Furthermore,

$$
\left\|\sigma-\sigma_{h}\right\|_{L}^{2}=\left\|\sigma-\operatorname{sym} \sigma_{h}\right\|_{L}^{2}+\left\|A s\left(\sigma_{h}\right)\right\|_{L}^{2}
$$

and $\left\|A s\left(\sigma_{h}\right)\right\|_{L}$ may simultaneously be regarded as an error contribution and as a (computable) contribution to an a posteriori error estimate. Hence the estimate (4.10) results in

$$
\begin{equation*}
\left\|\sigma-\sigma_{h}\right\|_{L} \approx\left\|\mathcal{R} e s_{N C}\right\|_{H^{*}}+\left\|\mathcal{R} e s_{H}\right\|_{H^{*}}+\left\|\operatorname{As}\left(\sigma_{h}\right)\right\|_{L} \tag{4.13}
\end{equation*}
$$

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